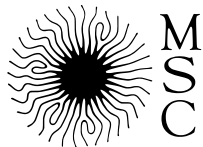


# Nonequilibrium signatures and phase transitions in active matter and beyond

PhD thesis of D. Martin supervised by J. Tailleur

Special thanks: C. Nardini

Collaborators: T. Arnoux de Pirey, D. Bartolo, H. Chaté, D. Geyer, Y. Kafri, M. Kardar, C. Nardini, J. O'byrne, A. Solon, F. Van Wijland



Laboratoire MSC  
université de Paris



October 8, 2021

# Microscopic Active Matter

- Particles exerting self-propulsion forces on their medium

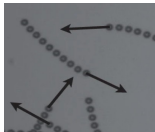
# Microscopic Active Matter

- Particles exerting self-propulsion forces on their medium

## Synthetic

- \* Physical mechanism

Quincke rollers



[Bricard et al, *Nature* 503 ]

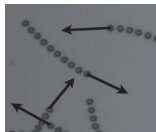
# Microscopic Active Matter

- Particles exerting self-propulsion forces on their medium

## Synthetic

- \* Physical mechanism

Quincke rollers

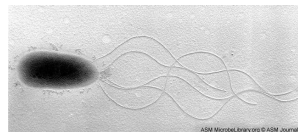


[Bricard et al, *Nature* 503 ]

## Living entities

- \* Biological mechanism

Bacterium



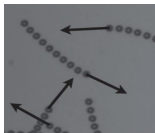
# Microscopic Active Matter

- Particles exerting self-propulsion forces on their medium

## Synthetic

- \* Physical mechanism

Quincke rollers

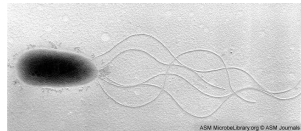


[Bricard et al, *Nature* 503 ]

## Living entities

- \* Biological mechanism

Bacterium



- Simplest theoretical models → non-Gaussian correlated fluctuations

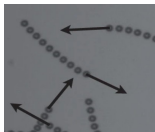
# Microscopic Active Matter

- Particles exerting self-propulsion forces on their medium

## Synthetic

- \* Physical mechanism

Quincke rollers

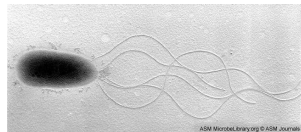


[Bricard et al, *Nature* 503 ]

## Living entities

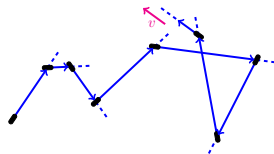
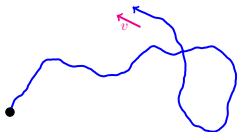
- \* Biological mechanism

Bacterium



- Simplest theoretical models  $\rightarrow$  non-Gaussian correlated fluctuations

ABP



RTP

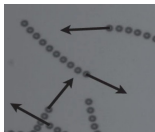
# Microscopic Active Matter

- Particles exerting self-propulsion forces on their medium

## Synthetic

- \* Physical mechanism

Quincke rollers

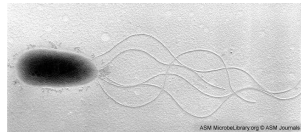


[Bricard et al, *Nature* 503]

## Living entities

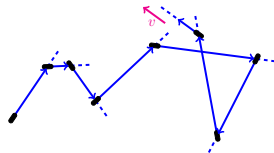
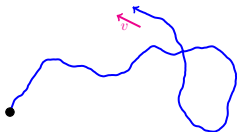
- \* Biological mechanism

Bacterium



- Simplest theoretical models  $\rightarrow$  non-Gaussian correlated fluctuations

ABP



RTP

- unusual fluctuations  $\rightarrow$  algebraic computations **challenging**

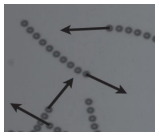
# Microscopic Active Matter

- Particles exerting self-propulsion forces on their medium

## Synthetic

- \* Physical mechanism

Quincke rollers

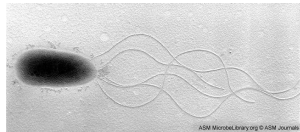


[Bricard et al, *Nature* 503]

## Living entities

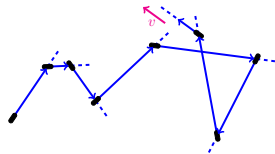
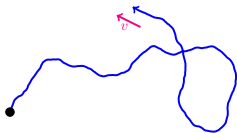
- \* Biological mechanism

Bacterium



- Simplest theoretical models  $\rightarrow$  non-Gaussian correlated fluctuations

ABP



RTP

- unusual fluctuations  $\rightarrow$  algebraic computations **challenging**
- steady-state distribution: **unknown**
- departure from equilibrium: **unquantified**



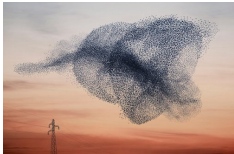
# Macroscopic Active Matter

- Active systems → ubiquitous in nature

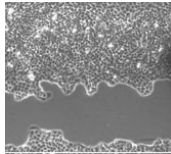
# Macroscopic Active Matter

- Active systems → ubiquitous in nature

## Biological



[Ballerini et al, *PNAS* 105, 2008]

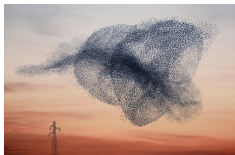


[Poujade et al, *PNAS* 104, 2007]

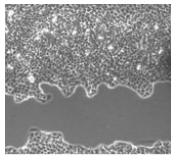
# Macroscopic Active Matter

- Active systems → ubiquitous in nature

## Biological



[Ballerini et al, *PNAS* 105, 2008]

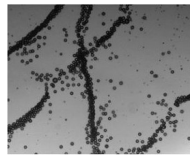


[Poujade et al, *PNAS* 104, 2007]

## Synthetic



[Geyer et al, *PRX* 9, 2019]

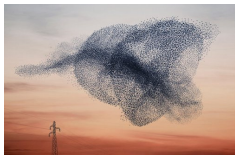


[Thutupalli et al, *PNAS* 115, 2018]

# Macroscopic Active Matter

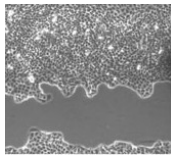
- Active systems → ubiquitous in nature

## Biological



[Ballerini et al, *PNAS* 105, 2008]

Bird flocks



[Poujade et al, *PNAS* 104, 2007]

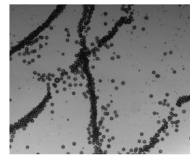
Healing tissue

## Synthetic



[Geyer et al, *PRX* 9, 2019]

Colloidal flocks



[Thutupalli et al, *PNAS* 115, 2018]

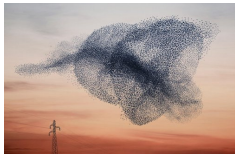
Lines of active droplets

- Self-organization emerges from collective dynamics → active phase transitions

# Macroscopic Active Matter

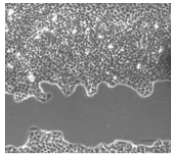
- Active systems → ubiquitous in nature

## Biological



[Ballerini et al, *PNAS* 105, 2008]

Bird flocks



[Poujade et al, *PNAS* 104, 2007]

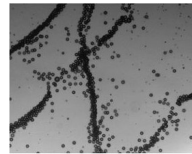
Healing tissue

## Synthetic



[Geyer et al, *PRX* 9, 2019]

Colloidal flocks



[Thutupalli et al, *PNAS* 115, 2018]

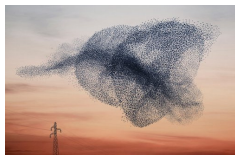
Lines of active droplets

- Self-organization emerges from collective dynamics → active phase transitions
- Controlling active phases → first step for engineering active materials

# Macroscopic Active Matter

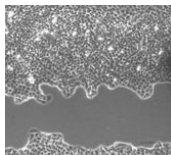
- Active systems → ubiquitous in nature

## Biological



[Ballerini et al, *PNAS* 105, 2008]

Bird flocks



[Poujade et al, *PNAS* 104, 2007]

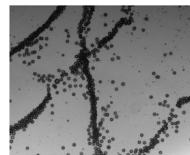
Healing tissue

## Synthetic



[Geyer et al, *PRX* 9, 2019]

Colloidal flocks



[Thutupalli et al, *PNAS* 115, 2018]

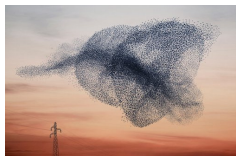
Lines of active droplets

- Self-organization emerges from collective dynamics → active phase transitions
- Controlling active phases → first step for engineering active materials
- Statistical physics: minimal ingredients

# Macroscopic Active Matter

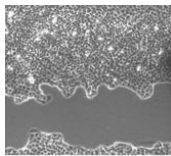
- Active systems → ubiquitous in nature

## Biological



[Ballerini et al, *PNAS* 105, 2008]

Bird flocks



[Poujade et al, *PNAS* 104, 2007]

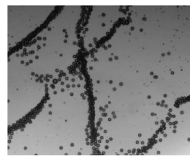
Healing tissue

## Synthetic



[Geyer et al, *PRX* 9, 2019]

Colloidal flocks



[Thutupalli et al, *PNAS* 115, 2018]

Lines of active droplets

- Self-organization emerges from collective dynamics → active phase transitions
- Controlling active phases → first step for engineering active materials
- Statistical physics: minimal ingredients

Motility-Induced Phase Separation (MIPS)  
repulsive forces

Flocking transition  
alignment

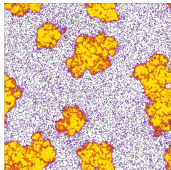
# Statistical physics of Active Matter: MIPS

- Phase separation at equilibrium: attractive forces vs thermal noise
  - ↳ Low  $T$  → cohesion wins: liquid-gas coexistence

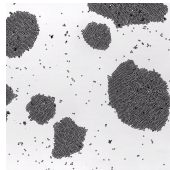


# Statistical physics of Active Matter: MIPS

- Phase separation at **equilibrium**: **attractive forces** vs **thermal noise**  
↳ Low  $T$  → cohesion wins: liquid-gas coexistence
- **Active particles**: repulsion triggers phase-separation (MIPS) [Tailleur et al, PRL 2008]



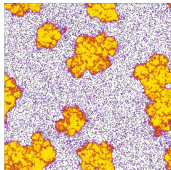
[Martin et al. PRE 2021]



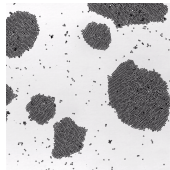
[Van Der Linden et al. PRL 2019]

# Statistical physics of Active Matter: MIPS

- Phase separation at **equilibrium**: **attractive forces** vs **thermal noise**  
↳ Low  $T$  → cohesion wins: liquid-gas coexistence
- **Active particles**: repulsion triggers phase-separation (MIPS) [Tailleur et al, PRL 2008]



[Martin et al. PRE 2021]

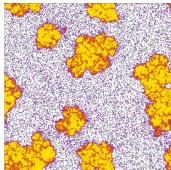


[Van Der Linden et al. PRL 2019]

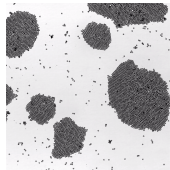
What is the **mechanism** behind MIPS ?

# Statistical physics of Active Matter: MIPS

- Phase separation at **equilibrium**: **attractive forces** vs **thermal noise**  
↳ Low  $T$  → cohesion wins: liquid-gas coexistence
- **Active particles**: repulsion triggers phase-separation (MIPS) [Tailleur et al, PRL 2008]



[Martin et al. PRE 2021]



[Van Der Linden et al. PRL 2019]

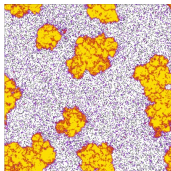
What is the **mechanism** behind MIPS ?

\* active particle **accumulate** in slow regions

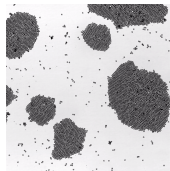
\* Repulsion **slows down** particles

# Statistical physics of Active Matter: MIPS

- Phase separation at **equilibrium**: **attractive forces** vs **thermal noise**  
↳ Low  $T$  → cohesion wins: liquid-gas coexistence
- **Active particles**: repulsion triggers phase-separation (MIPS) [Tailleur et al, PRL 2008]



[Martin et al. PRE 2021]



[Van Der Linden et al. PRL 2019]

What is the **mechanism** behind MIPS ?

\* active particle **accumulate** in slow regions

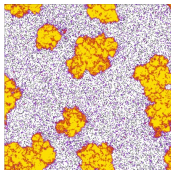
\* Repulsion **slows down** particles



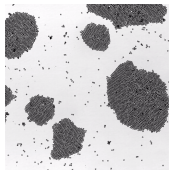
formation of **dense clusters**

# Statistical physics of Active Matter: MIPS

- Phase separation at **equilibrium**: **attractive forces** vs **thermal noise**  
↳ Low  $T$  → cohesion wins: liquid-gas coexistence
- **Active particles**: repulsion triggers phase-separation (MIPS) [Tailleur et al, PRL 2008]



[Martin et al. PRE 2021]



[Van Der Linden et al. PRL 2019]

What is the **mechanism** behind MIPS ?

\* active particle **accumulate** in slow regions

\* Repulsion **slows down** particles



formation of **dense clusters**

- MIPS in **self-propelled spheres**: starts to be **understood**
- MIPS for **generic interactions**: **more complex**

# Statistical physics of Active Matter: flocking

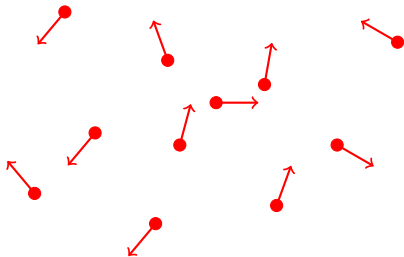
- Modelling flocking transition → aligning interactions

# Statistical physics of Active Matter: flocking

- Modelling flocking transition → aligning interactions  
    ↳ Ferromagnetic alignment → minimal active model: Vicsek Model

# Statistical physics of Active Matter: flocking

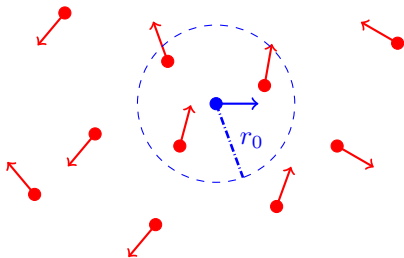
- Modelling flocking transition  $\rightarrow$  aligning interactions  
     $\hookrightarrow$  Ferromagnetic alignment  $\rightarrow$  minimal active model: Vicsek Model
- Point-like flying spins at speed  $v$





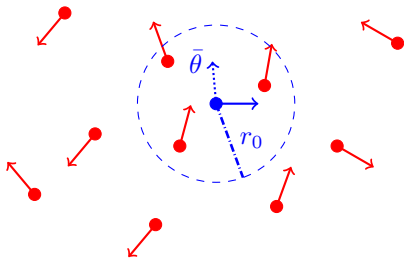
# Statistical physics of Active Matter: flocking

- Modelling flocking transition  $\rightarrow$  aligning interactions  
     $\hookrightarrow$  Ferromagnetic alignment  $\rightarrow$  minimal active model: Vicsek Model
- Point-like flying spins at speed  $v$



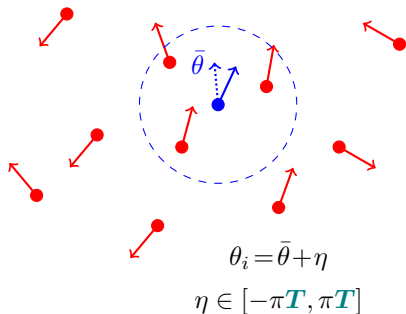
# Statistical physics of Active Matter: flocking

- Modelling flocking transition  $\rightarrow$  aligning interactions  
     $\hookrightarrow$  Ferromagnetic alignment  $\rightarrow$  minimal active model: Vicsek Model
- Point-like flying spins at speed  $v$



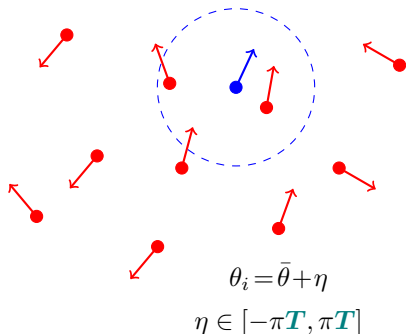
# Statistical physics of Active Matter: flocking

- Modelling flocking transition  $\rightarrow$  aligning interactions  
     $\hookrightarrow$  Ferromagnetic alignment  $\rightarrow$  minimal active model: Vicsek Model
- Point-like flying spins at speed  $v$



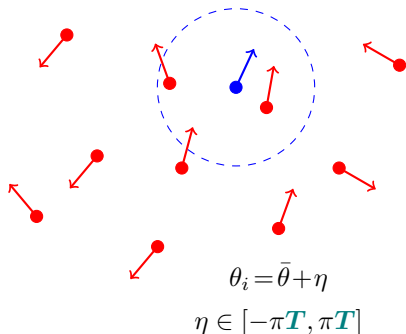
# Statistical physics of Active Matter: flocking

- Modelling flocking transition  $\rightarrow$  aligning interactions  
     $\hookrightarrow$  Ferromagnetic alignment  $\rightarrow$  minimal active model: Vicsek Model
- Point-like flying spins at speed  $v$



# Statistical physics of Active Matter: flocking

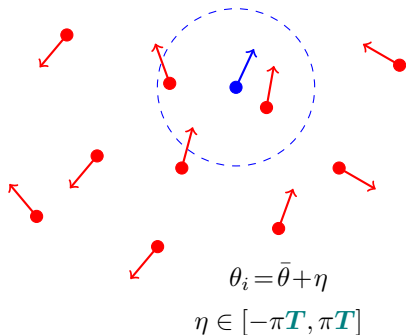
- Modelling flocking transition  $\rightarrow$  aligning interactions  
     $\hookrightarrow$  Ferromagnetic alignment  $\rightarrow$  minimal active model: Vicsek Model
- Point-like flying spins at speed  $v$



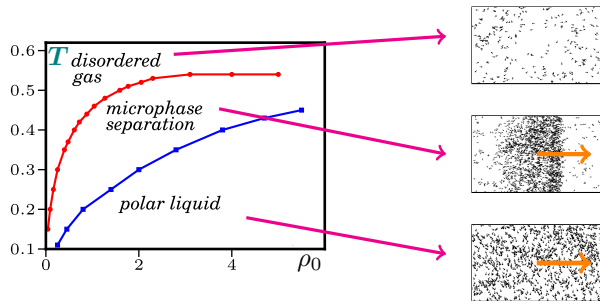
# Statistical physics of Active Matter: flocking

- Modelling flocking transition  $\rightarrow$  aligning interactions  
 $\hookrightarrow$  Ferromagnetic alignment  $\rightarrow$  minimal active model: Vicsek Model

- Point-like flying spins at speed  $v$



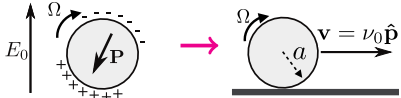
- Phase diagram [Solon et al, PRL 114, 2015]



# Statistical physics of Active Matter: flocking

- Vicsek Model  $\rightarrow$  relevant for experiments

Quincke rollers



[Bricard et al, *Nature* 503 ]

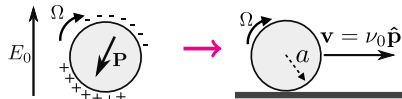
$\Rightarrow$   
 $\Rightarrow$   
Pack them



# Statistical physics of Active Matter: flocking

- Vicsek Model  $\rightarrow$  relevant for experiments

Quincke rollers



[Bricard et al, *Nature* 503 ]

$\Rightarrow$   
 $\Rightarrow$   
Pack them



- hydrodynamic + electrostatic interactions

$\rightarrow$   
[Bricard et al, *Nature* 503 ]

effective alignment



# Statistical physics of Active Matter: flocking

- Vicsek Model  $\rightarrow$  relevant for experiments

Quincke rollers



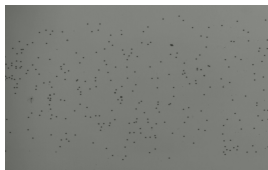
[Bricard et al, *Nature* 503 ]

$\Rightarrow$   
 $\Rightarrow$   
Pack them

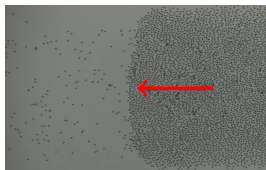


- hydrodynamic + electrostatic interactions  $\rightarrow$  effective alignment  
[Bricard et al, *Nature* 503 ]

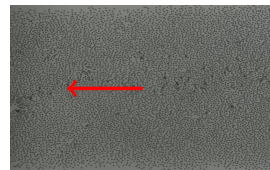
- Experimental phase diagram similar to the Vicsek Model



Disordered gas



Polar bands

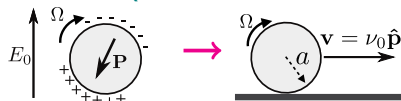


Ordered flock

# Statistical physics of Active Matter: flocking

- Vicsek Model  $\rightarrow$  relevant for experiments

Quincke rollers



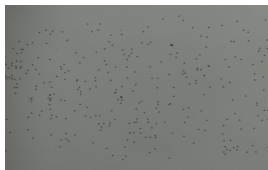
[Bricard et al, *Nature* 503 ]

$\Rightarrow$   
 $\Rightarrow$   
Pack them

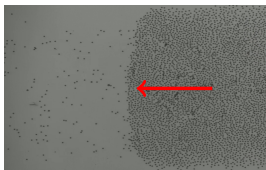


- hydrodynamic + electrostatic interactions  $\rightarrow$  effective alignment  
[Bricard et al, *Nature* 503 ]

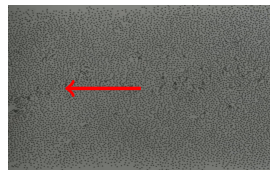
- Experimental phase diagram similar to the Vicsek Model



Disordered gas



Polar bands



Ordered flock

- Emergence of flocks in the Vicsek Model: starts to be understood

# Statistical physics of Active Matter: flocking

- Vicsek Model  $\rightarrow$  relevant for experiments

Quincke rollers



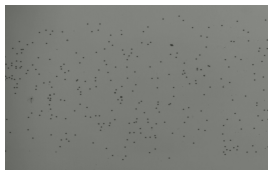
[Bricard et al, *Nature* 503 ]

$\Rightarrow$   
 $\Rightarrow$   
Pack them

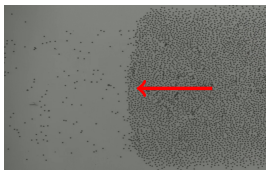


- hydrodynamic + electrostatic interactions  $\rightarrow$  effective alignment  
[Bricard et al, *Nature* 503 ]

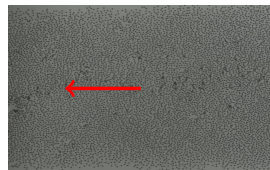
- Experimental phase diagram similar to the Vicsek Model



Disordered gas



Polar bands



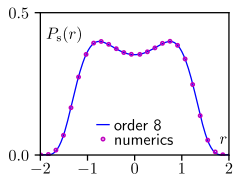
Ordered flock

- Emergence of flocks in the Vicsek Model: starts to be understood

What lies beyond for more complex systems ?

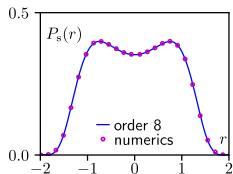
# The four axes of this thesis

## Exact results for a single active particle

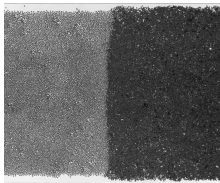


# The four axes of this thesis

## Exact results for a single active particle

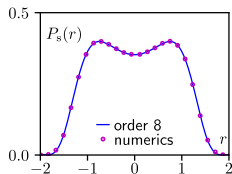


## MIPS in dense polar flocks

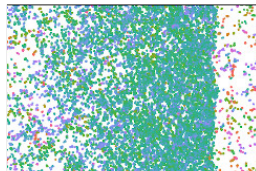


# The four axes of this thesis

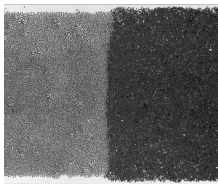
## Exact results for a single active particle



## Fluctuation-induced first-order flocking

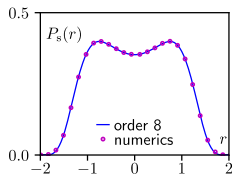


## MIPS in dense polar flocks

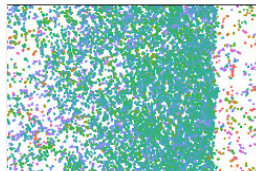


# The four axes of this thesis

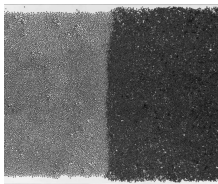
## Exact results for a single active particle



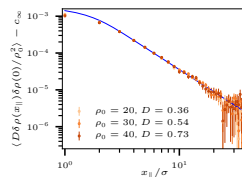
## Fluctuation-induced first-order flocking



## MIPS in dense polar flocks



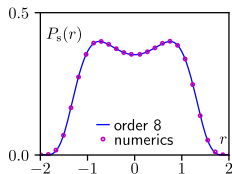
## Anisotropy-induced long-ranged correlations



# The four axes of this thesis

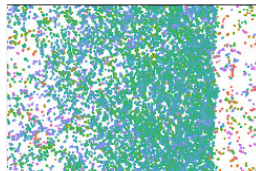
## Part I

Exact results for a single active particle



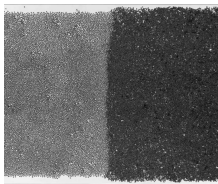
## Part III

Fluctuation-induced first-order flocking

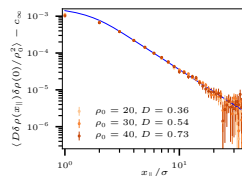


## Part II

MIPS in dense polar flocks



Anisotropy-induced long-ranged correlations





# Microscopic Active Matter: exact approaches

- Simplest model  $\rightarrow$  Active Ornstein-Uhlenbeck Particles (AOUPs) [Fodor et al, *PRL* 117, 2016]

$$\dot{x} = -\partial_x \phi + v, \quad \tau \dot{v} = -v + \sqrt{2D} \eta \quad \text{with} \quad \langle v(t)v(t') \rangle = \frac{D}{\tau} e^{-|t-t'|/\tau}$$

# Microscopic Active Matter: exact approaches

- Simplest model  $\rightarrow$  Active Ornstein-Uhlenbeck Particles (AOUPs) [Fodor et al, *PRL* 117, 2016]

$$\dot{x} = -\partial_x \phi + v, \quad \tau \dot{v} = -v + \sqrt{2D} \eta \quad \text{with} \quad \langle v(t)v(t') \rangle = \frac{D}{\tau} e^{-|t-t'|/\tau}$$

- $\tau \sim 0 \iff$  Equilibrium, temperature  $D$   $\left\{ \begin{array}{l} \text{Boltzmann distribution} \rightarrow P_s(x) = e^{-\frac{\phi}{D}} \\ \text{No steady-state current} \rightarrow J = \langle \dot{x} \rangle = 0 \end{array} \right.$

# Microscopic Active Matter: exact approaches

- Simplest model  $\rightarrow$  Active Ornstein-Uhlenbeck Particles (AOUPs) [Fodor et al, *PRL* 117, 2016]

$$\dot{x} = -\partial_x \phi + v, \quad \tau \dot{v} = -v + \sqrt{2D} \eta \quad \text{with} \quad \langle v(t)v(t') \rangle = \frac{D}{\tau} e^{-|t-t'|/\tau}$$

- $\tau \sim 0 \Leftrightarrow$  Equilibrium, temperature  $D$   $\left\{ \begin{array}{l} \text{Boltzmann distribution} \rightarrow P_s(x) = e^{-\frac{\phi}{D}} \\ \text{No steady-state current} \rightarrow J = \langle \dot{x} \rangle = 0 \end{array} \right.$
- $\tau \neq 0 \Leftrightarrow$  out-of-equilibrium  $\left\{ \begin{array}{l} \text{deviation from Boltzmann} \rightarrow P_s(x) = e^{-\frac{\phi}{D}} + \tau \dots + \tau^2 \dots + \dots \\ \text{Nonzero ratchet current} \rightarrow J = \tau^2 \dots + \dots \end{array} \right.$

# Microscopic Active Matter: exact approaches

- Simplest model  $\rightarrow$  Active Ornstein-Uhlenbeck Particles (AOUPs) [Fodor et al, *PRL* 117, 2016]

$$\dot{x} = -\partial_x \phi + v, \quad \tau \dot{v} = -v + \sqrt{2D} \eta \quad \text{with} \quad \langle v(t)v(t') \rangle = \frac{D}{\tau} e^{-|t-t'|/\tau}$$

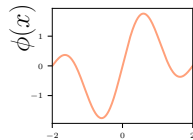
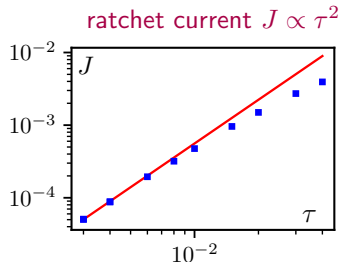
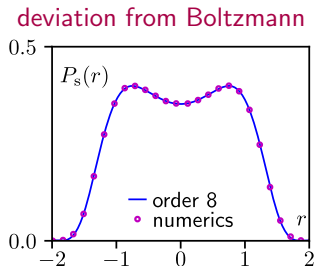
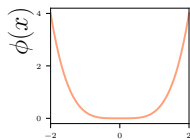
- $\tau \sim 0 \Leftrightarrow$  Equilibrium, temperature  $D$   $\left\{ \begin{array}{l} \text{Boltzmann distribution} \rightarrow P_s(x) = e^{-\frac{\phi}{D}} \\ \text{No steady-state current} \rightarrow J = \langle \dot{x} \rangle = 0 \end{array} \right.$
- $\tau \neq 0 \Leftrightarrow$  out-of-equilibrium  $\left\{ \begin{array}{l} \text{deviation from Boltzmann} \rightarrow P_s(x) = e^{-\frac{\phi}{D}} + \tau \dots + \tau^2 \dots + \dots \\ \text{Nonzero ratchet current} \rightarrow J = \tau^2 \dots + \dots \end{array} \right.$

# Microscopic Active Matter: exact approaches

- Simplest model  $\rightarrow$  Active Ornstein-Uhlenbeck Particles (AOUPs) [Fodor et al, *PRL* 117, 2016]

$$\dot{x} = -\partial_x \phi + v, \quad \tau \dot{v} = -v + \sqrt{2D} \eta \quad \text{with} \quad \langle v(t)v(t') \rangle = \frac{D}{\tau} e^{-|t-t'|/\tau}$$

- $\tau \sim 0 \Leftrightarrow$  Equilibrium, temperature  $D$ 
  - Boltzmann distribution  $\rightarrow P_s(x) = e^{-\frac{\phi}{D}}$
  - No steady-state current  $\rightarrow J = \langle \dot{x} \rangle = 0$
- $\tau \neq 0 \Leftrightarrow$  out-of-equilibrium
  - deviation from Boltzmann  $\rightarrow P_s(x) = e^{-\frac{\phi}{D}} + \tau \dots + \tau^2 \dots + \dots$
  - Nonzero ratchet current  $\rightarrow J = \tau^2 \dots + \dots$



# Microscopic Active matter: exact approaches

- Active Matter models → neglects thermal fluctuations

## Microscopic Active matter: exact approaches

- Active Matter models → neglects thermal fluctuations
  - ↳ interplay between active and passive noises rarely studied

# Microscopic Active matter: exact approaches

- Active Matter models  $\rightarrow$  neglects thermal fluctuations  
 $\hookrightarrow$  interplay between active and passive noises rarely studied
- Minimal model: the AOUP with thermal noise

$$\dot{x} = -\partial_x \phi + v + \sqrt{2T} \eta_1, \quad \tau \dot{v} = -v + \sqrt{2D} \eta_2$$



# Microscopic Active matter: exact approaches

- Active Matter models  $\rightarrow$  neglects thermal fluctuations  
 $\hookrightarrow$  interplay between active and passive noises rarely studied
- Minimal model: the AOUP with thermal noise

$$\dot{x} = -\partial_x \phi + v + \sqrt{2T} \eta_1, \quad \tau \dot{v} = -v + \sqrt{2D} \eta_2$$

$\hookrightarrow \tau = 0 \iff$  Equilibrium at temperature  $T + D$

## Microscopic Active matter: exact approaches

- Active Matter models  $\rightarrow$  neglects thermal fluctuations  
 $\hookrightarrow$  interplay between active and passive noises rarely studied

- Minimal model: the AOUP with thermal noise

$$\dot{x} = -\partial_x \phi + v + \sqrt{2T} \eta_1, \quad \tau \dot{v} = -v + \sqrt{2D} \eta_2$$

$\hookrightarrow \tau = 0 \iff$  Equilibrium at temperature  $T + D$

- $\tau \neq 0$ : effect of  $T$  on nonequilibrium signatures ?

# Microscopic Active matter: exact approaches

- Active Matter models  $\rightarrow$  neglects thermal fluctuations  
 $\hookrightarrow$  interplay between active and passive noises rarely studied

- Minimal model: the AOUP with thermal noise

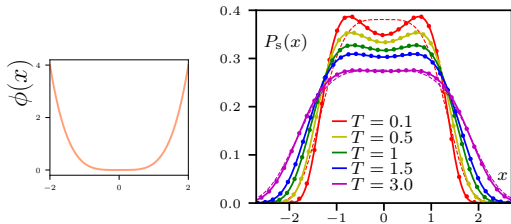
$$\dot{x} = -\partial_x \phi + v + \sqrt{2T} \eta_1, \quad \tau \dot{v} = -v + \sqrt{2D} \eta_2$$

$\hookrightarrow \tau = 0 \iff$  Equilibrium at temperature  $T + D$

- $\tau \neq 0$ : effect of  $T$  on nonequilibrium signatures ?

$T$  seems to wash out activity...

deviation from Boltzmann



# Microscopic Active matter: exact approaches

- Active Matter models  $\rightarrow$  neglects thermal fluctuations  
 $\hookrightarrow$  interplay between active and passive noises rarely studied

- Minimal model: the AOUP with thermal noise

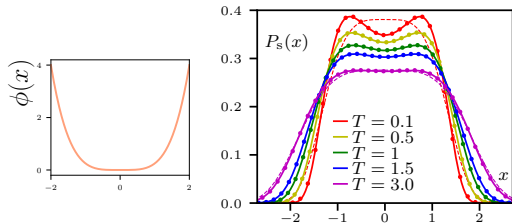
$$\dot{x} = -\partial_x \phi + v + \sqrt{2T} \eta_1, \quad \tau \dot{v} = -v + \sqrt{2D} \eta_2$$

$\hookrightarrow \tau = 0 \iff$  Equilibrium at temperature  $T + D$

- $\tau \neq 0$ : effect of  $T$  on nonequilibrium signatures ?

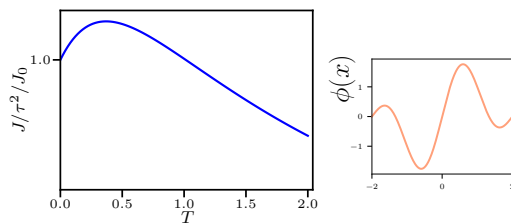
$T$  seems to wash out activity...

deviation from Boltzmann



...but enhances it

ratchet current  $J$



# Microscopic Active matter: exact approaches

- Active Matter models  $\rightarrow$  neglects thermal fluctuations  
 $\hookrightarrow$  interplay between active and passive noises rarely studied

- Minimal model: the AOUP with thermal noise

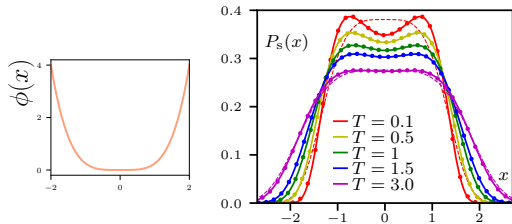
$$\dot{x} = -\partial_x \phi + v + \sqrt{2T} \eta_1, \quad \tau \dot{v} = -v + \sqrt{2D} \eta_2$$

$\hookrightarrow \tau = 0 \iff$  Equilibrium at temperature  $T + D$

- $\tau \neq 0$ : effect of  $T$  on nonequilibrium signatures ?

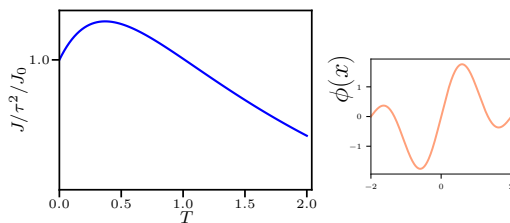
$T$  seems to wash out activity...

deviation from Boltzmann



...but enhances it

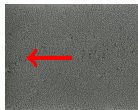
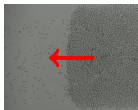
ratchet current  $J$



- Non trivial interplay between active and passive noises

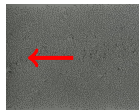
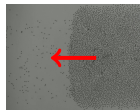
## Part II: Motility-induced solidification in roller flocks

Low density: Vicsek

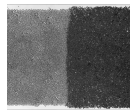


## Part II: Motility-induced solidification in roller flocks

Low density: Vicsek



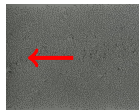
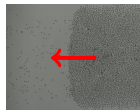
High density: beyond Vicsek



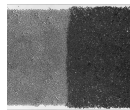
New phase transition  
active solidification

## Part II: Motility-induced solidification in roller flocks

Low density: Vicsek



High density: beyond Vicsek



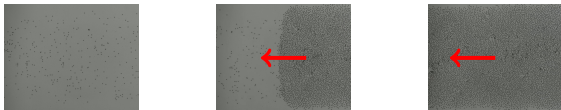
New phase transition  
active solidification

- What is happening ?

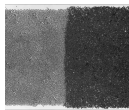


## Part II: Motility-induced solidification in roller flocks

Low density: Vicsek

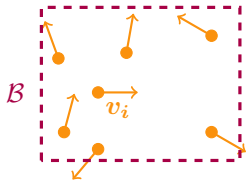


High density: beyond Vicsek



New phase transition  
active solidification

- What is happening ?

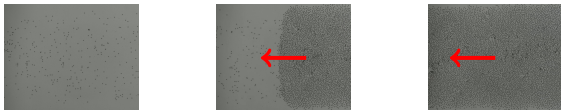


$$W = \left| \frac{1}{N} \sum_{i \in \mathcal{B}} v_i \right|$$

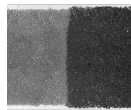
$$v = \frac{1}{N} \sum_{i \in \mathcal{B}} |v_i|$$

## Part II: Motility-induced solidification in roller flocks

Low density: Vicsek

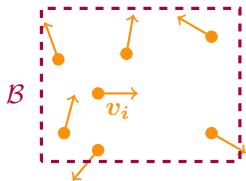


High density: beyond Vicsek



New phase transition  
active solidification

- What is happening ?

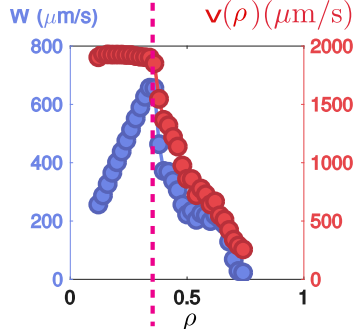


$$W = \left| \frac{1}{N} \sum_{i \in \mathcal{B}} v_i \right|$$

$$v = \frac{1}{N} \sum_{i \in \mathcal{B}} |v_i|$$

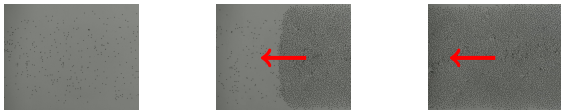
Vicsek physics

Beyond Vicsek

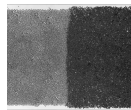


## Part II: Motility-induced solidification in roller flocks

Low density: Vicsek

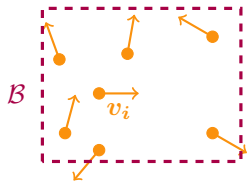


High density: beyond Vicsek



New phase transition  
active solidification

- What is happening ?



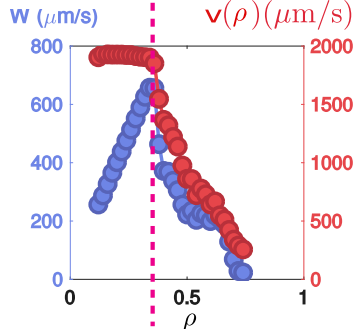
$$W = \left| \frac{1}{N} \sum_{i \in B} v_i \right|$$

$$v = \frac{1}{N} \sum_{i \in B} |v_i|$$

⇒ Is it MIPS at play in a flock ?

Vicsek physics

Beyond Vicsek



# Phenomenological hydrodynamics

- Standard hydrodynamic of the Vicsek model for  $\rho = \langle \sum_i \delta(r - r_i) \rangle$  and  $W = \langle \sum_i v_i \delta(r - r_i) \rangle$

$$\partial_t \rho = D_\rho \partial_{xx} \rho - \partial_x W \quad (1)$$

$$\partial_t W + \lambda W \partial_x W = D_W \partial_{xx} W - \partial_x (v\rho) + \alpha W - a_4 W^3 \quad (2)$$

# Phenomenological hydrodynamics

- Standard **hydrodynamic** of the Vicsek model for  $\rho = \langle \sum_i \delta(r - r_i) \rangle$  and  $W = \langle \sum_i v_i \delta(r - r_i) \rangle$

$$\partial_t \rho = D_\rho \partial_{xx} \rho - \partial_x W \quad (1)$$

$$\partial_t W + \lambda W \partial_x W = D_W \partial_{xx} W - \partial_x (v \rho) + \alpha W - a_4 W^3 \quad (2)$$

- Adapt it to high-density experimental features
  - Rollers' velocity drop:  $v \rightarrow v(\rho)$
  - Rollers lose orientational order:  $\alpha \rightarrow \alpha(\rho)$

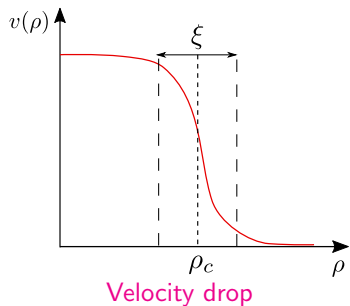
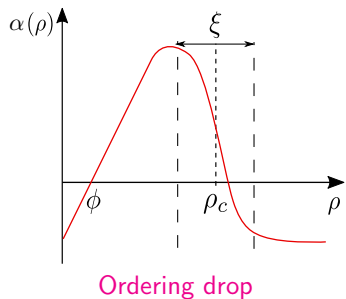
# Phenomenological hydrodynamics

- Standard hydrodynamic of the Vicsek model for  $\rho = \langle \sum_i \delta(r - r_i) \rangle$  and  $W = \langle \sum_i v_i \delta(r - r_i) \rangle$

$$\partial_t \rho = D_\rho \partial_{xx} \rho - \partial_x W \quad (1)$$

$$\partial_t W + \lambda W \partial_x W = D_W \partial_{xx} W - \partial_x (v\rho) + \alpha W - a_4 W^3 \quad (2)$$

- Adapt it to high-density experimental features
  - Rollers' velocity drop:  $v \rightarrow v(\rho)$
  - Rollers lose orientational order:  $\alpha \rightarrow \alpha(\rho)$



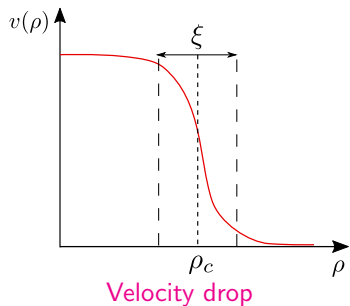
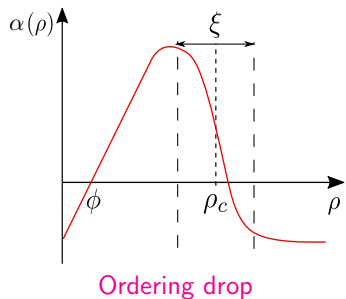
# Phenomenological hydrodynamics

- Standard hydrodynamic of the Vicsek model for  $\rho = \langle \sum_i \delta(r - r_i) \rangle$  and  $W = \langle \sum_i v_i \delta(r - r_i) \rangle$

$$\partial_t \rho = D_\rho \partial_{xx} \rho - \partial_x W \quad (1)$$

$$\partial_t W + \lambda W \partial_x W = D_W \partial_{xx} W - \partial_x (v\rho) + \alpha W - a_4 W^3 \quad (2)$$

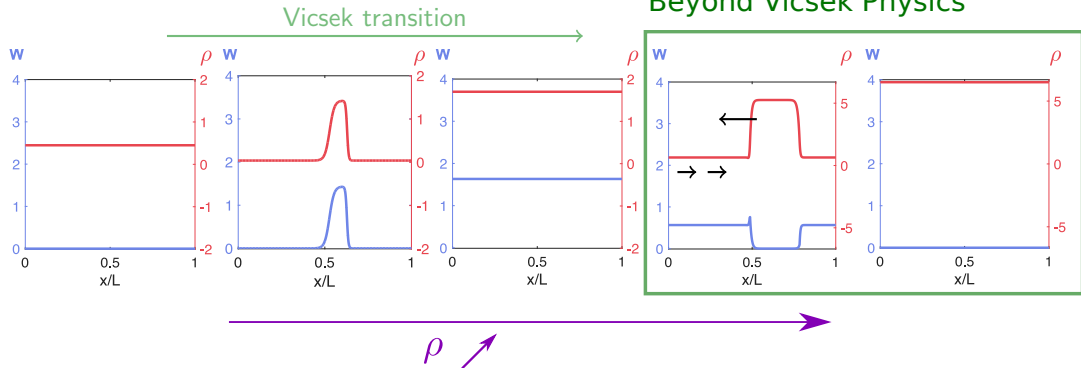
- Adapt it to high-density experimental features
  - Rollers' velocity drop:  $v \rightarrow v(\rho)$
  - Rollers lose orientational order:  $\alpha \rightarrow \alpha(\rho)$



Postulated phenomenologically, could be rigorously derived

# Phenomenological hydrodynamics

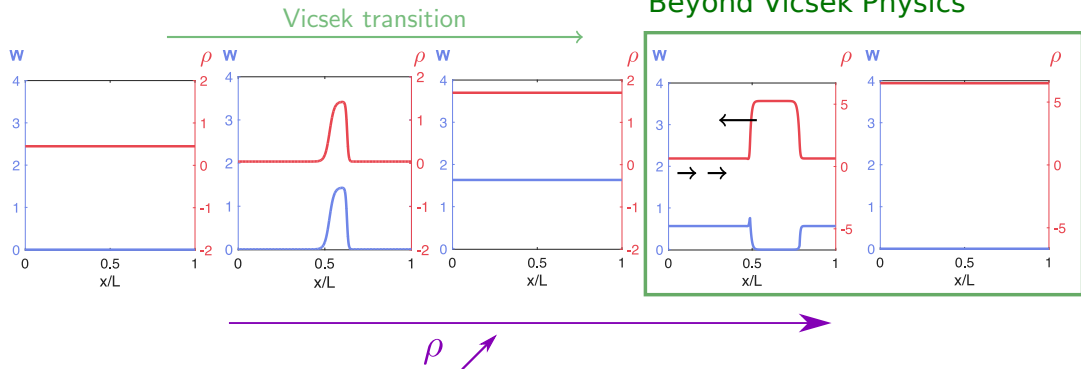
- New phase transition at high density





# Phenomenological hydrodynamics

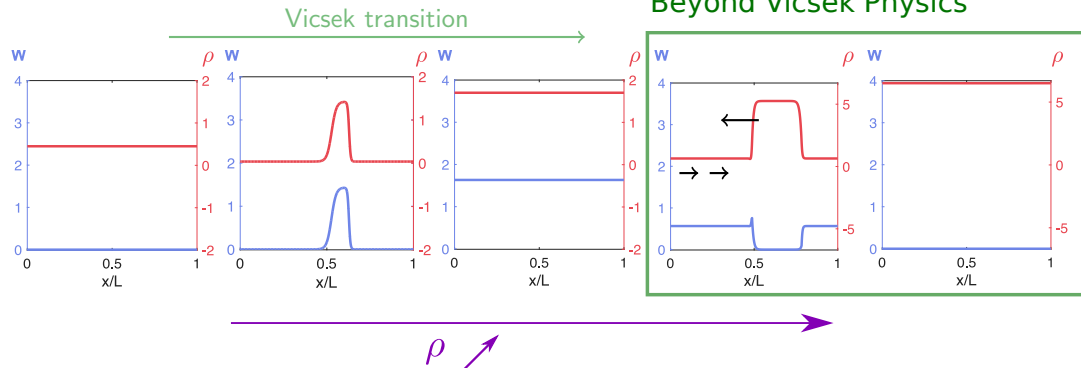
- New phase transition at high density



- MIPS-like transition: linear instability, lever rule, hysteresis loops, coarsening dynamics.

# Phenomenological hydrodynamics

- New phase transition at high density



## Beyond Vicsek Physics

- MIPS-like transition: linear instability, lever rule, hysteresis loops, coarsening dynamics.
- Movie •

# Quincke rollers experiments

Vicsek transition

Gas



Vicsek Bands

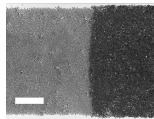


Polar Liquid

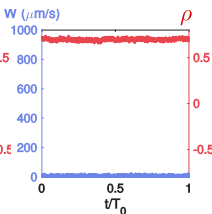
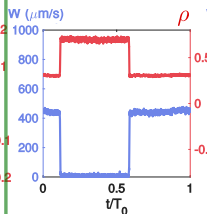
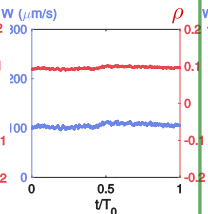
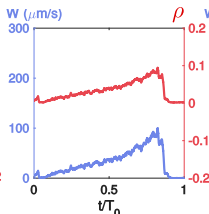
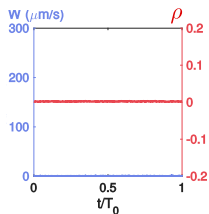
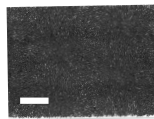


Beyond Vicsek Physics

Solid Jam



Active Solid



$\rho$  ↑

# Quincke rollers experiments

Vicsek transition

Gas



Vicsek Bands

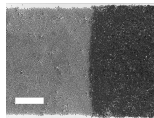


Polar Liquid

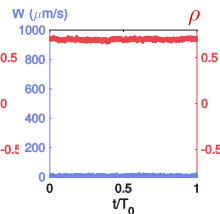
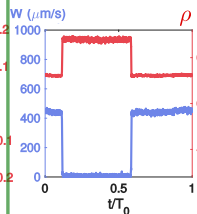
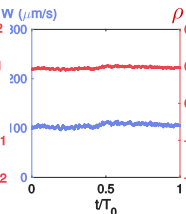
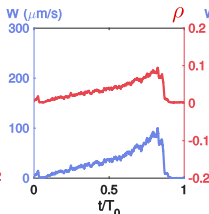
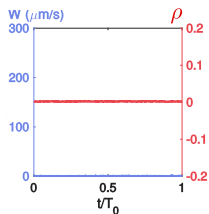
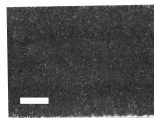


Beyond Vicsek Physics

Solid Jam



Active Solid



$\rho$  ↗

- Striking similarities of the phase diagram

# Quincke rollers experiments

Vicsek transition

Gas



Vicsek Bands

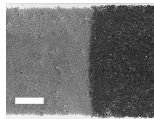


Polar Liquid

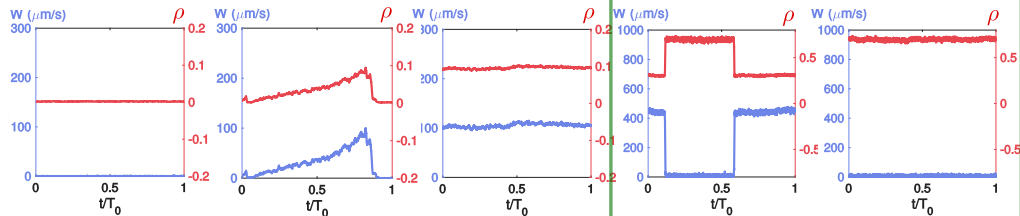
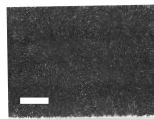


Beyond Vicsek Physics

Solid Jam



Active Solid



$\rho$  ↗

- Striking similarities of the phase diagram

Movie •

- New phase transition at high  $\rho$  in roller flocks  $\rightarrow$  active solidification

## Part II: conclusion

- New phase transition at high  $\rho$  in roller flocks  $\rightarrow$  active solidification
- Described by MIPS occurring in a polar liquid

## Part II: conclusion

- New phase transition at high  $\rho$  in roller flocks  $\rightarrow$  active solidification
- Described by MIPS occurring in a polar liquid
- speed reduction + dense flocks of active units = active solidification



## Part II: conclusion

- New phase transition at high  $\rho$  in roller flocks  $\rightarrow$  active solidification
- Described by MIPS occurring in a polar liquid
- speed reduction + dense flocks of active units = active solidification
- Agreement of phenomenological hydrodynamics and experiments beyond the phase diagram
  - $\hookrightarrow$  Lever rule
  - $\hookrightarrow$  Hysteresis loops

## Part III: Fluctuation-induced phase separation in models of collective motion

- Emergence of collective motion:  $1^{st}$  or  $2^{nd}$  order ?
  - Practical importance: existence of a band phase

## Part III: Fluctuation-induced phase separation in models of collective motion

- Emergence of collective motion:  $1^{st}$  or  $2^{nd}$  order ?  
→ Practical importance: existence of a band phase
- Long standing debate {
  - numerics: strong finite-size effects
  - analytics: mean-field approximations

## Part III: Fluctuation-induced phase separation in models of collective motion

- Emergence of collective motion:  $1^{st}$  or  $2^{nd}$  order ?  
→ Practical importance: existence of a band phase

- Long standing debate {  
  numerics: strong finite-size effects  
  analytics: mean-field approximations

Recipe:    flocking models → hydrodynamics → linear stability → order of the transition

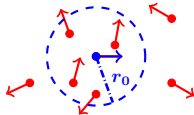
## Part III: Fluctuation-induced phase separation in models of collective motion

- Emergence of collective motion:  $1^{st}$  or  $2^{nd}$  order ?  
→ Practical importance: existence of a band phase
- Long standing debate {
  - numerics: strong finite-size effects
  - analytics: mean-field approximations

**Recipe:** flocking models → hydrodynamics → linear stability → order of the transition

### First order / coexistence

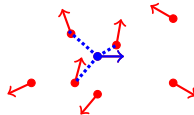
- Vicsek: metric alignment



[Grégoire et al, PRL 92, 2004]

### Second order / continuous

- Vicsek: topological alignment



[Chou et al, PRE 86, 2012]

- Active Ising Model: hydrodynamics

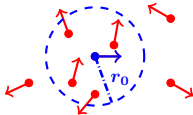
## Part III: Fluctuation-induced phase separation in models of collective motion

- Emergence of collective motion:  $1^{st}$  or  $2^{nd}$  order ?  
→ Practical importance: existence of a band phase
- Long standing debate {
  - numerics: strong finite-size effects
  - analytics: mean-field approximations

**Recipe:** flocking models → hydrodynamics → linear stability → order of the transition

### First order / coexistence

- Vicsek: metric alignment

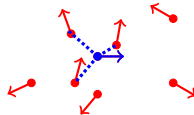


[Grégoire et al, PRL 92, 2004]

- Active Ising Model: numerics

### Second order / continuous

- Vicsek: topological alignment



[Chou et al, PRE 86, 2012]

- Active Ising Model: hydrodynamics

## Part III: Fluctuation-induced phase separation in models of collective motion

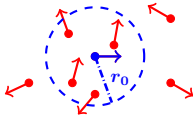
- Emergence of collective motion:  $1^{st}$  or  $2^{nd}$  order ?  
→ Practical importance: existence of a band phase

- Long standing debate {
  - numerics: strong finite-size effects
  - analytics: mean-field approximations

**Recipe:** flocking models → hydrodynamics → linear stability → order of the transition

### First order / coexistence

- Vicsek: metric alignment

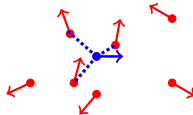


[Grégoire et al, PRL 92, 2004]

- Active Ising Model: numerics

### Second order / continuous

- Vicsek: topological alignment



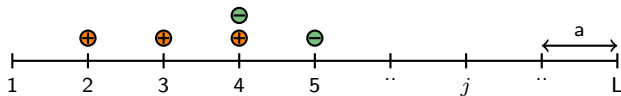
[Chou et al, PRE 86, 2012]

- Active Ising Model: hydrodynamics

→ However, microscopic simulations disagree: why? Is it generic?

# The Active Ising Model [Solon et al, PRL 111, 2013]

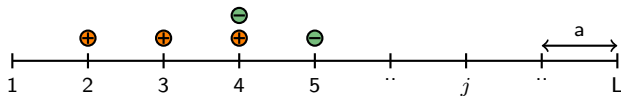
- specific system but generic results
- a simplified spin model for collective motion





# The Active Ising Model [Solon et al, PRL 111, 2013]

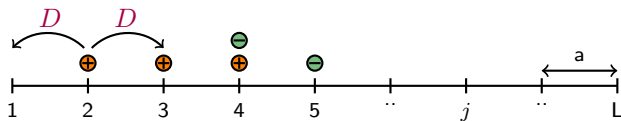
- specific system but generic results
- a simplified spin model for collective motion



- Density  $\rho_j \simeq (n_j^+ + n_j^-)/a$       Magnetization  $m_j \simeq (n_j^+ - n_j^-)/a$

# The Active Ising Model [Solon et al, PRL 111, 2013]

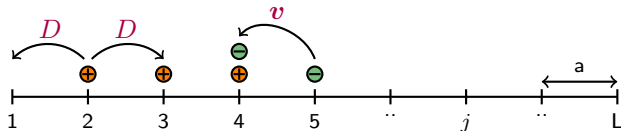
- specific system but generic results
- a simplified spin model for collective motion



- Density  $\rho_j \simeq (n_j^+ + n_j^-)/a$       Magnetization  $m_j \simeq (n_j^+ - n_j^-)/a$
- Isotropic diffusion with rate  $D$

# The Active Ising Model [Solon et al, PRL 111, 2013]

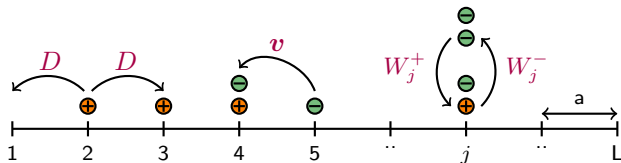
- specific system but generic results
- a simplified spin model for collective motion



- Density  $\rho_j \simeq (n_j^+ + n_j^-)/a$       Magnetization  $m_j \simeq (n_j^+ - n_j^-)/a$
- Isotropic diffusion with rate  $D$
- Active jumps with rate  $v$  :  $\ominus$  jumps left,  $\oplus$  jumps right.

# The Active Ising Model [Solon et al, PRL 111, 2013]

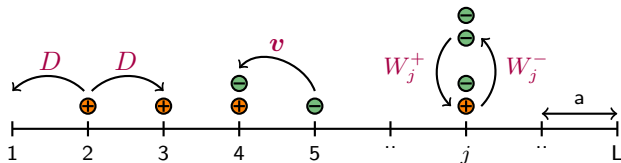
- specific system but generic results
- a simplified spin model for collective motion



- Density  $\rho_j \simeq (n_j^+ + n_j^-)/a$       Magnetization  $m_j \simeq (n_j^+ - n_j^-)/a$
- Isotropic diffusion with rate  $D$
- Active jumps with rate  $v$  :  $\ominus$  jumps left,  $\oplus$  jumps right.
- Spins align with rate  $W_j^\pm = \exp\left(\pm\beta\frac{m_j}{\rho_j}\right) \Leftrightarrow$  fully connected Ising Model on site  $j$

# The Active Ising Model [Solon et al, PRL 111, 2013]

- specific system but generic results
- a simplified spin model for collective motion



- Density  $\rho_j \simeq (n_j^+ + n_j^-)/a$       Magnetization  $m_j \simeq (n_j^+ - n_j^-)/a$
- Isotropic diffusion with rate  $D$
- Active jumps with rate  $v$  :  $\ominus$  jumps left,  $\oplus$  jumps right.
- Spins align with rate  $W_j^\pm = \exp\left(\pm\beta\frac{m_j}{\rho_j}\right) \Leftrightarrow$  fully connected Ising Model on site  $j$

Master equation + Mean-field approximation



Hydrodynamics for  $\rho$  and  $m$

# The Active Ising Model [Solon et al, PRL 111, 2013]

$$\partial_t \rho = D_\rho \partial_{xx} \rho - v \partial_x m$$

$$\partial_t m = D_m \partial_{xx} m - v \partial_x \rho + 2\rho \sinh\left(\beta \frac{m}{\rho}\right) - 2m \cosh\left(\beta \frac{m}{\rho}\right)$$

# The Active Ising Model [Solon et al, PRL 111, 2013]

$$\partial_t \rho = D_\rho \partial_{xx} \rho - v \partial_x m$$

$$\partial_t m = D_m \partial_{xx} m - v \partial_x \rho + 2\rho \sinh\left(\beta \frac{m}{\rho}\right) - 2m \cosh\left(\beta \frac{m}{\rho}\right)$$

- Landau expansion:  $2\rho \sinh\left(\beta \frac{m}{\rho}\right) - 2m \cosh\left(\beta \frac{m}{\rho}\right) = \alpha m - \gamma \frac{m^3}{\rho^2} + \mathcal{O}(m^5)$

# The Active Ising Model [Solon et al, PRL 111, 2013]

$$\partial_t \rho = D_\rho \partial_{xx} \rho - v \partial_x m$$

$$\partial_t m = D_m \partial_{xx} m - v \partial_x \rho + \alpha m - \gamma \frac{m^3}{\rho^2}$$

- Landau expansion:  $2\rho \sinh\left(\beta \frac{m}{\rho}\right) - 2m \cosh\left(\beta \frac{m}{\rho}\right) = \alpha m - \gamma \frac{m^3}{\rho^2} + \mathcal{O}(m^5)$
- Nonequilibrium model  $\rightarrow$  Landau terms  $\mathcal{F}_{MF} = \alpha m - \gamma \frac{m^3}{\rho^2}$  do not derive from a free energy



# The Active Ising Model [Solon et al, PRL 111, 2013]

$$\partial_t \rho = D_\rho \partial_{xx} \rho - v \partial_x m$$

$$\partial_t m = D_m \partial_{xx} m - v \partial_x \rho + \alpha m - \gamma \frac{m^3}{\rho^2}$$

- Landau expansion:  $2\rho \sinh\left(\beta \frac{m}{\rho}\right) - 2m \cosh\left(\beta \frac{m}{\rho}\right) = \alpha m - \gamma \frac{m^3}{\rho^2} + \mathcal{O}(m^5)$
- Nonequilibrium model  $\rightarrow$  Landau terms  $\mathcal{F}_{MF} = \alpha m - \gamma \frac{m^3}{\rho^2}$  do not derive from a free energy
- Linear stability of homogeneous profiles  
 $\hookrightarrow \alpha = \text{constant} \rightarrow$  ordered solution linearly stable at onset

# The Active Ising Model [Solon et al, PRL 111, 2013]

$$\partial_t \rho = D_\rho \partial_{xx} \rho - v \partial_x m$$

$$\partial_t m = D_m \partial_{xx} m - v \partial_x \rho + \alpha m - \gamma \frac{m^3}{\rho^2}$$

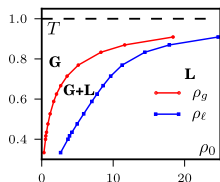
- Landau expansion:  $2\rho \sinh\left(\beta \frac{m}{\rho}\right) - 2m \cosh\left(\beta \frac{m}{\rho}\right) = \alpha m - \gamma \frac{m^3}{\rho^2} + \mathcal{O}(m^5)$
- Nonequilibrium model  $\rightarrow$  Landau terms  $\mathcal{F}_{MF} = \alpha m - \gamma \frac{m^3}{\rho^2}$  do not derive from a free energy
- Linear stability of homogeneous profiles
  - $\hookrightarrow \alpha = \text{constant} \rightarrow$  ordered solution linearly stable at onset
  - $\hookrightarrow$  continuous emergence of flocking

# The Active Ising Model [Solon et al, PRL 111, 2013]

$$\partial_t \rho = D_\rho \partial_{xx} \rho - v \partial_x m$$

$$\partial_t m = D_m \partial_{xx} m - v \partial_x \rho + \alpha m - \gamma \frac{m^3}{\rho^2}$$

- Landau expansion:  $2\rho \sinh\left(\beta \frac{m}{\rho}\right) - 2m \cosh\left(\beta \frac{m}{\rho}\right) = \alpha m - \gamma \frac{m^3}{\rho^2} + \mathcal{O}(m^5)$
- Nonequilibrium model  $\rightarrow$  Landau terms  $\mathcal{F}_{MF} = \alpha m - \gamma \frac{m^3}{\rho^2}$  do not derive from a free energy
- Linear stability of homogeneous profiles
  - $\rightarrow \alpha = \text{constant} \rightarrow$  ordered solution linearly stable at onset
  - $\rightarrow$  continuous emergence of flocking



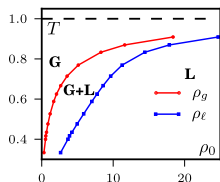
$\Rightarrow$  At odds with microscopic simulations

# The Active Ising Model [Solon et al, PRL 111, 2013]

$$\partial_t \rho = D_\rho \partial_{xx} \rho - v \partial_x m$$

$$\partial_t m = D_m \partial_{xx} m - v \partial_x \rho + \alpha m - \gamma \frac{m^3}{\rho^2}$$

- Landau expansion:  $2\rho \sinh\left(\beta \frac{m}{\rho}\right) - 2m \cosh\left(\beta \frac{m}{\rho}\right) = \alpha m - \gamma \frac{m^3}{\rho^2} + \mathcal{O}(m^5)$
- Nonequilibrium model  $\rightarrow$  Landau terms  $\mathcal{F}_{MF} = \alpha m - \gamma \frac{m^3}{\rho^2}$  do not derive from a free energy
- Linear stability of homogeneous profiles
  - $\rightarrow \alpha = \text{constant} \rightarrow$  ordered solution linearly stable at onset
  - $\rightarrow$  continuous emergence of flocking



$\Rightarrow$  At odds with microscopic simulations

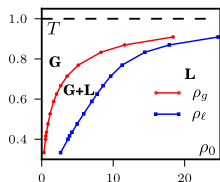
- What did mean-field miss ?

# The Active Ising Model [Solon et al, PRL 111, 2013]

$$\partial_t \rho = D_\rho \partial_{xx} \rho - v \partial_x m$$

$$\partial_t m = D_m \partial_{xx} m - v \partial_x \rho + \alpha m - \gamma \frac{m^3}{\rho^2} + \sqrt{2\sigma\rho} \eta$$

- Landau expansion:  $2\rho \sinh\left(\beta \frac{m}{\rho}\right) - 2m \cosh\left(\beta \frac{m}{\rho}\right) = \alpha m - \gamma \frac{m^3}{\rho^2} + \mathcal{O}(m^5)$
- Nonequilibrium model  $\rightarrow$  Landau terms  $\mathcal{F}_{MF} = \alpha m - \gamma \frac{m^3}{\rho^2}$  do not derive from a free energy
- Linear stability of homogeneous profiles
  - $\rightarrow \alpha = \text{constant} \rightarrow$  ordered solution linearly stable at onset
  - $\rightarrow$  continuous emergence of flocking



$\Rightarrow$  At odds with microscopic simulations

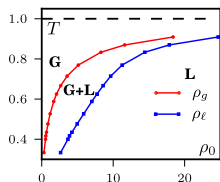
- What did mean-field miss? **Fluctuations**

# The Active Ising Model [Solon et al, PRL 111, 2013]

$$\partial_t \rho = D_\rho \partial_{xx} \rho - v \partial_x m$$

$$\partial_t m = D_m \partial_{xx} m - v \partial_x \rho + \alpha m - \gamma \frac{m^3}{\rho^2} + \sqrt{2\sigma\rho} \eta$$

- Landau expansion:  $2\rho \sinh\left(\beta \frac{m}{\rho}\right) - 2m \cosh\left(\beta \frac{m}{\rho}\right) = \alpha m - \gamma \frac{m^3}{\rho^2} + \mathcal{O}(m^5)$
- Nonequilibrium model  $\rightarrow$  Landau terms  $\mathcal{F}_{MF} = \alpha m - \gamma \frac{m^3}{\rho^2}$  do not derive from a free energy
- Linear stability of homogeneous profiles
  - $\rightarrow \alpha = \text{constant} \rightarrow$  ordered solution linearly stable at onset
  - $\rightarrow$  continuous emergence of flocking



$\Rightarrow$  At odds with microscopic simulations

- What did mean-field miss? **Fluctuations**
  - $\rightarrow$  renormalize the Landau terms in the dynamics of  $\langle \rho \rangle$  and  $\langle m \rangle$   $\mathcal{F} \rightarrow \mathcal{F}_{MF} + \Delta \mathcal{F}$

## Quasi-linear renormalization

- Fluctuations  $\rho = \rho_0 + \delta\rho$  and  $m = m_0 + \delta m$  around homogeneous mean field

## Quasi-linear renormalization

- Fluctuations  $\rho = \rho_0 + \delta\rho$  and  $m = m_0 + \delta m$  around homogeneous mean field  
    ↪ Landau terms  $\mathcal{F}$  perturbed



# Quasi-linear renormalization

- Fluctuations  $\rho = \rho_0 + \delta\rho$  and  $m = m_0 + \delta m$  around homogeneous mean field

↪ Landau terms  $\mathcal{F}$  perturbed

$$\Delta\mathcal{F} = \frac{\partial\mathcal{F}}{\partial m}\langle\delta m\rangle + \frac{\partial\mathcal{F}}{\partial\rho}\langle\delta\rho\rangle + \frac{\partial^2\mathcal{F}}{\partial m\partial\rho}\langle\delta\rho\delta m\rangle + \frac{1}{2}\frac{\partial^2\mathcal{F}}{\partial^2 m}\langle\delta m^2\rangle + \frac{1}{2}\frac{\partial^2\mathcal{F}}{\partial^2\rho}\langle\delta\rho^2\rangle$$

# Quasi-linear renormalization

- Fluctuations  $\rho = \rho_0 + \delta\rho$  and  $m = m_0 + \delta m$  around homogeneous mean field

↪ Landau terms  $\mathcal{F}$  perturbed

$$\Delta\mathcal{F} = \frac{\partial\mathcal{F}}{\partial m}\langle\delta m\rangle + \frac{\partial\mathcal{F}}{\partial\rho}\langle\delta\rho\rangle + \frac{\partial^2\mathcal{F}}{\partial m\partial\rho}\langle\delta\rho\delta m\rangle + \frac{1}{2}\frac{\partial^2\mathcal{F}}{\partial^2 m}\langle\delta m^2\rangle + \frac{1}{2}\frac{\partial^2\mathcal{F}}{\partial^2\rho}\langle\delta\rho^2\rangle$$

↪ Dynamics of  $\delta\rho$  and  $\delta m$  at linear order ?

# Quasi-linear renormalization

- Fluctuations  $\rho = \rho_0 + \delta\rho$  and  $m = m_0 + \delta m$  around homogeneous mean field

↪ Landau terms  $\mathcal{F}$  perturbed

$$\Delta\mathcal{F} = \frac{\partial\mathcal{F}}{\partial m}\langle\delta m\rangle + \frac{\partial\mathcal{F}}{\partial\rho}\langle\delta\rho\rangle + \frac{\partial^2\mathcal{F}}{\partial m\partial\rho}\langle\delta\rho\delta m\rangle + \frac{1}{2}\frac{\partial^2\mathcal{F}}{\partial^2 m}\langle\delta m^2\rangle + \frac{1}{2}\frac{\partial^2\mathcal{F}}{\partial^2\rho}\langle\delta\rho^2\rangle$$

↪ Dynamics of  $\delta\rho$  and  $\delta m$  at linear order ?

$$\partial_t \begin{pmatrix} \delta\rho_q \\ \delta m_q \end{pmatrix} = \begin{pmatrix} M_{11}^q & M_{12}^q \\ M_{21}^q & M_{22}^q \end{pmatrix} \begin{pmatrix} \delta\rho_q \\ \delta m_q \end{pmatrix} + \begin{pmatrix} 0 \\ \sqrt{2\sigma\rho_0}\eta_q \end{pmatrix}$$

# Quasi-linear renormalization

- Fluctuations  $\rho = \rho_0 + \delta\rho$  and  $m = m_0 + \delta m$  around homogeneous mean field

↪ Landau terms  $\mathcal{F}$  perturbed

$$\Delta\mathcal{F} = \frac{\partial\mathcal{F}}{\partial m}\langle\delta m\rangle + \frac{\partial\mathcal{F}}{\partial\rho}\langle\delta\rho\rangle + \frac{\partial^2\mathcal{F}}{\partial m\partial\rho}\langle\delta\rho\delta m\rangle + \frac{1}{2}\frac{\partial^2\mathcal{F}}{\partial^2 m}\langle\delta m^2\rangle + \frac{1}{2}\frac{\partial^2\mathcal{F}}{\partial^2\rho}\langle\delta\rho^2\rangle$$

↪ Dynamics of  $\delta\rho$  and  $\delta m$  at linear order ?

$$\partial_t \begin{pmatrix} \delta\rho_q \\ \delta m_q \end{pmatrix} = \begin{pmatrix} M_{11}^q & M_{12}^q \\ M_{21}^q & M_{22}^q \end{pmatrix} \begin{pmatrix} \delta\rho_q \\ \delta m_q \end{pmatrix} + \begin{pmatrix} 0 \\ \sqrt{2\sigma\rho_0}\eta_q \end{pmatrix}$$



steady-state correlators

# Quasi-linear renormalization

- Fluctuations  $\rho = \rho_0 + \delta\rho$  and  $m = m_0 + \delta m$  around homogeneous mean field

↪ Landau terms  $\mathcal{F}$  perturbed

$$\Delta\mathcal{F} = \frac{\partial\mathcal{F}}{\partial m}\langle\delta m\rangle + \frac{\partial\mathcal{F}}{\partial\rho}\langle\delta\rho\rangle + \frac{\partial^2\mathcal{F}}{\partial m\partial\rho}\langle\delta\rho\delta m\rangle + \frac{1}{2}\frac{\partial^2\mathcal{F}}{\partial^2 m}\langle\delta m^2\rangle + \frac{1}{2}\frac{\partial^2\mathcal{F}}{\partial^2\rho}\langle\delta\rho^2\rangle$$

↪ Dynamics of  $\delta\rho$  and  $\delta m$  at linear order ?

$$\partial_t \begin{pmatrix} \delta\rho_q \\ \delta m_q \end{pmatrix} = \begin{pmatrix} M_{11}^q & M_{12}^q \\ M_{21}^q & M_{22}^q \end{pmatrix} \begin{pmatrix} \delta\rho_q \\ \delta m_q \end{pmatrix} + \begin{pmatrix} 0 \\ \sqrt{2\sigma\rho_0} \eta_q \end{pmatrix}$$



steady-state correlators

$$\langle\delta m\rangle = 0$$

$$\langle\delta\rho\rangle = 0$$

$$\langle\delta\rho\delta m\rangle = 0 + \mathcal{O}(m_0)$$

$$\langle\delta m^2\rangle = \sigma\rho_0 \frac{v^2 \sqrt{\frac{2\alpha}{D} + \alpha\sqrt{v^2 + \alpha D}}}{8\alpha v^2 + 4\alpha^2 D} + \mathcal{O}(m_0)$$

$$\langle\delta\rho^2\rangle = \sigma\rho_0 \frac{v^2 \left( \sqrt{\frac{2\alpha}{D} - \frac{\alpha}{\sqrt{v^2 + \alpha D}}} \right)}{4\alpha(\alpha D + 2v^2)} + \mathcal{O}(m_0)$$

# Quasi-linear renormalization

- Fluctuations  $\rho = \rho_0 + \delta\rho$  and  $m = m_0 + \delta m$  around homogeneous mean field

↪ Landau terms  $\mathcal{F}$  perturbed

$$\Delta\mathcal{F} = \frac{\partial\mathcal{F}}{\partial m} \langle \delta m \rangle + \frac{\partial\mathcal{F}}{\partial \rho} \langle \delta \rho \rangle + \frac{\partial^2\mathcal{F}}{\partial m \partial \rho} \langle \delta \rho \delta m \rangle + \frac{1}{2} \frac{\partial^2\mathcal{F}}{\partial^2 m} \langle \delta m^2 \rangle + \frac{1}{2} \frac{\partial^2\mathcal{F}}{\partial^2 \rho} \langle \delta \rho^2 \rangle$$

↪ Dynamics of  $\delta\rho$  and  $\delta m$  at linear order ?

$$\partial_t \begin{pmatrix} \delta\rho_q \\ \delta m_q \end{pmatrix} = \begin{pmatrix} M_{11}^q & M_{12}^q \\ M_{21}^q & M_{22}^q \end{pmatrix} \begin{pmatrix} \delta\rho_q \\ \delta m_q \end{pmatrix} + \begin{pmatrix} 0 \\ \sqrt{2\sigma\rho_0} \eta_q \end{pmatrix}$$



steady-state correlators

$$\langle \delta m \rangle = 0$$

$$\langle \delta \rho \rangle = 0$$

$$\langle \delta \rho \delta m \rangle = 0 + \mathcal{O}(m_0)$$

$$\langle \delta m^2 \rangle = \sigma\rho_0 \frac{v^2 \sqrt{\frac{2\alpha}{D} + \alpha\sqrt{v^2 + \alpha D}}}{8\alpha v^2 + 4\alpha^2 D} + \mathcal{O}(m_0)$$

$$\langle \delta \rho^2 \rangle = \sigma\rho_0 \frac{v^2 \left( \sqrt{\frac{2\alpha}{D} - \frac{\alpha}{\sqrt{v^2 + \alpha D}}} \right)}{4\alpha(\alpha D + 2v^2)} + \mathcal{O}(m_0)$$

# Quasi-linear renormalization

- Landau terms  $\mathcal{F} \rightarrow \mathcal{F} + \Delta\mathcal{F}$

# Quasi-linear renormalization

- Landau terms  $\mathcal{F} \rightarrow \mathcal{F} + \Delta\mathcal{F}$

$$\alpha \rightarrow \alpha + \frac{\sigma\gamma}{\rho v} f\left(\frac{\alpha D}{v^2}\right) + \mathcal{O}(\sigma^2) \Rightarrow \alpha \text{ density dependent}$$



# Quasi-linear renormalization

- Landau terms  $\mathcal{F} \rightarrow \mathcal{F} + \Delta\mathcal{F}$

$$\alpha \rightarrow \alpha + \frac{\sigma\gamma}{\rho v} f\left(\frac{\alpha D}{v^2}\right) + \mathcal{O}(\sigma^2) \Rightarrow \alpha \text{ density dependent}$$

- Linear stability analysis for  $\alpha(\rho)$

↳ Homogeneous **polar** and **disordered** profiles **unstable** at onset

# Quasi-linear renormalization

- Landau terms  $\mathcal{F} \rightarrow \mathcal{F} + \Delta\mathcal{F}$

$$\alpha \longrightarrow \alpha + \frac{\sigma\gamma}{\rho v} f\left(\frac{\alpha D}{v^2}\right) + \mathcal{O}(\sigma^2) \quad \Rightarrow \alpha \text{ density dependent}$$

- Linear stability analysis for  $\alpha(\rho)$

↳ Homogeneous polar and disordered profiles unstable at onset

↳ Discontinuous emergence of flocking

# Quasi-linear renormalization

- Landau terms  $\mathcal{F} \rightarrow \mathcal{F} + \Delta\mathcal{F}$

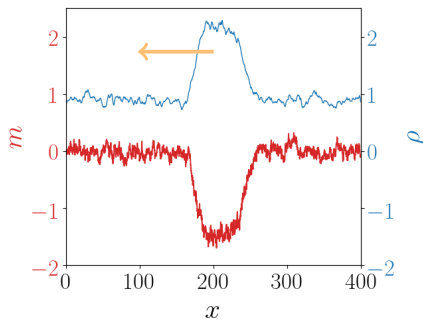
$$\alpha \rightarrow \alpha + \frac{\sigma\gamma}{\rho v} f\left(\frac{\alpha D}{v^2}\right) + \mathcal{O}(\sigma^2) \Rightarrow \alpha \text{ density dependent}$$

- Linear stability analysis for  $\alpha(\rho)$

↳ Homogeneous polar and disordered profiles unstable at onset

↳ Discontinuous emergence of flocking

- Simulations of the stochastic PDE



# Fluctuations makes the transition discontinuous

- Summary of the mechanism in Active Ising Model

# Fluctuations makes the transition discontinuous

- Summary of the mechanism in Active Ising Model

MF hydrodynamics with deceptive continuous transition

# Fluctuations makes the transition discontinuous

- Summary of the mechanism in Active Ising Model

MF hydrodynamics with deceptive continuous transition



# Fluctuations makes the transition discontinuous

- Summary of the mechanism in Active Ising Model

MF hydrodynamics with deceptive continuous transition



Makes linear Landau term  $\propto$  density dependent

# Fluctuations makes the transition discontinuous

- Summary of the mechanism in Active Ising Model

MF hydrodynamics with deceptive continuous transition

⇓ + fluctuations

Makes linear Landau term  $\alpha$  density dependent

⇓  $\alpha(\rho)$



# Fluctuations makes the transition discontinuous

- Summary of the mechanism in Active Ising Model

MF hydrodynamics with deceptive continuous transition

⇓ + fluctuations

Makes linear Landau term  $\alpha$  density dependent

⇓  $\alpha(\rho)$

Makes the transition discontinuous

# Fluctuations makes the transition discontinuous

- Summary of the mechanism in Active Ising Model

MF hydrodynamics with deceptive continuous transition

↓ ↓ + fluctuations

Makes linear Landau term  $\alpha$  density dependent

↓ ↓  $\alpha(\rho)$

Makes the transition discontinuous

- Metric models: Fluctuation-induced first-order transition

# Fluctuations makes the transition discontinuous

- Summary of the mechanism in Active Ising Model

MF hydrodynamics with deceptive continuous transition

⇓ + fluctuations

Makes linear Landau term  $\alpha$  density dependent

⇓  $\alpha(\rho)$

Makes the transition discontinuous

- Metric models: Fluctuation-induced first-order transition  
↳ discontinuous transition with coexistence

# Fluctuations makes the transition discontinuous

- Summary of the mechanism in Active Ising Model

MF hydrodynamics with deceptive continuous transition

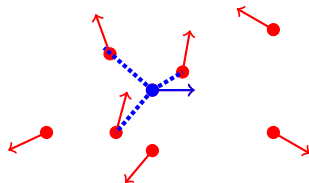


Makes linear Landau term  $\alpha$  density dependent



Makes the transition discontinuous

- Metric models: Fluctuation-induced first-order transition  
↳ discontinuous transition with coexistence



What about 'topological' or 'metric-free' models ?

# Topological models: a specific transition

- Visual or biological cues → 'metric-free' or 'topological' alignment

# Topological models: a specific transition

- Visual or biological cues → 'metric-free' or 'topological' alignment

$k$ -nearest neighbours



## Interaction ruling animal collective behavior depends on topological rather than metric distance: Evidence from a field study

M. Ballerini<sup>1\*</sup>, N. Cabibbo<sup>2,5</sup>, R. Candelier<sup>1,5</sup>, A. Cavagna<sup>1,6\*</sup>, E. Cisbani<sup>7</sup>, I. Giardina<sup>8</sup>, V. Lecomte<sup>1,10,11</sup>, A. Orlandi<sup>9</sup>, G. Parisi<sup>12,13,14\*</sup>, A. Procaccini<sup>15</sup>, and M. Viale<sup>1,15</sup>, and V. Zdravkovic<sup>6</sup>

<sup>1</sup>Centre for Statistical Mechanics and Complexity (SMC), Consiglio Nazionale delle Ricerche-Istituto Nazionale per la Fisica della Materia, <sup>2</sup>Dipartimento di Fisica, and <sup>3</sup>Sezione Istituto Nazionale di Fisica Nucleare, Università di Roma "La Sapienza," Piazzale Aldo Moro 2, 00185 Roma, Italy; <sup>4</sup>Istituto Superiore di Sanità, viale Regina Elena 299, 00161 Roma, Italy; <sup>5</sup>Istituto dei Sistemi Complessi (ISC), Consiglio Nazionale delle Ricerche, via dei Taurini 19, 00185 Roma, Italy; and <sup>6</sup>Laboratoire Matière et Systèmes Complexes, (Centre National de la Recherche Scientifique-Unité Mixte de Recherche 7057), Université Paris VII, 10 rue Alice Demon et Léonie Duquet, 75205 Paris Cedex 13, France

PHYSICAL REVIEW X **6**, 021011 (2016)

---

## Motility-Driven Glass and Jamming Transitions in Biological Tissues

Depeng Bi,<sup>1,3</sup> Xingbo Yang,<sup>1,4</sup> M. Cristina Marchetti,<sup>1,2</sup> and M. Lisa Manning<sup>1,2</sup>

Cell motion inside dense tissues governs many biological processes, including embryonic development and cancer metastasis, and recent experiments suggest that these tissues exhibit collective glassy behavior. To make quantitative predictions about glass transitions in tissues, we study a self-propelled Voronoi model that simultaneously captures polarized cell motility and multibody cell-cell interactions in a confluent

Voronoi neighbours

# Topological models: a specific transition

- Visual or biological cues → 'metric-free' or 'topological' alignment

$k$ -nearest neighbours



## Interaction ruling animal collective behavior depends on topological rather than metric distance: Evidence from a field study

M. Ballerini<sup>1</sup>, N. Cabibbo<sup>1,5</sup>, R. Candelier<sup>1,5</sup>, A. Cavagna<sup>1,6\*</sup>, E. Cisbani<sup>1</sup>, I. Giardina<sup>6</sup>, V. Lecomte<sup>1,7,8,9</sup>, A. Orlandi<sup>1</sup>, G. Parisi<sup>1,10,11\*</sup>, A. Procaccini<sup>1</sup>, and M. Viale<sup>1,12</sup>, and V. Zdravkovic<sup>1</sup>

<sup>1</sup>Centre for Statistical Mechanics and Complexity (SMC), Consiglio Nazionale delle Ricerche-Istituto Nazionale per la Fisica della Materia, <sup>2</sup>Dipartimento di Fisica, and <sup>3</sup>Sezione Istituto Nazionale di Fisica Nucleare, Università di Roma "La Sapienza," Piazzale Aldo Moro 2, 00185 Roma, Italy; <sup>4</sup>Istituto Superiore di Sanità, viale Regina Elena 299, 00161 Roma, Italy; <sup>5</sup>Istituto dei Sistemi Complessi (ISC), Consiglio Nazionale delle Ricerche, via dei Taurini 19, 00185 Roma, Italy; and <sup>6</sup>Laboratoire Matière et Systèmes Complexes, (Centre National de la Recherche Scientifique-Unité Mixte de Recherche 7057), Université Paris VII, 10 rue Alice Demon et Léonie Duquet, 75205 Paris Cedex 13, France

PHYSICAL REVIEW X **6**, 021011 (2016)

## Motility-Driven Glass and Jamming Transitions in Biological Tissues

Dapeng Bi,<sup>1,3</sup> Xingbo Yang,<sup>1,4</sup> M. Cristina Marchetti,<sup>1,2</sup> and M. Lisa Manning<sup>1,2</sup>

Cell motion inside dense tissues governs many biological processes, including embryonic development and cancer metastasis, and recent experiments suggest that these tissues exhibit collective glassy behavior. To make quantitative predictions about glass transitions in tissues, we study a self-propelled Voronoi model that simultaneously captures polarized cell motility and multibody cell-cell interactions in a confluent

- topological models:  $1^{st}$  or  $2^{nd}$  order flocking transition?

# Topological models: a specific transition

- Visual or biological cues → 'metric-free' or 'topological' alignment

$k$ -nearest neighbours



## Interaction ruling animal collective behavior depends on topological rather than metric distance: Evidence from a field study

M. Ballerini<sup>1</sup>, N. Cabibbo<sup>2,5</sup>, R. Candelier<sup>1,5</sup>, A. Cavagna<sup>1,6\*</sup>, E. Cisbani<sup>1</sup>, I. Giardina<sup>4</sup>, V. Lecomte<sup>1,7,8,9</sup>, A. Orlandi<sup>1</sup>, G. Parisi<sup>10,11,12\*</sup>, A. Procaccini<sup>11</sup>, and M. Viale<sup>13,5</sup>, and V. Zdravkovic<sup>1</sup>

<sup>1</sup>Centre for Statistical Mechanics and Complexity (SMC), Consiglio Nazionale delle Ricerche-Istituto Nazionale per la Fisica della Materia, <sup>2</sup>Dipartimento di Fisica, and <sup>3</sup>Sezione Istituto Nazionale di Fisica Nucleare, Università di Roma "La Sapienza," Piazzale Aldo Moro 2, 00185 Roma, Italy; <sup>4</sup>Istituto Superiore di Sanità, viale Regina Elena 299, 00161 Roma, Italy; <sup>5</sup>Istituto dei Sistemi Complessi (ISC), Consiglio Nazionale delle Ricerche, via dei Taurini 19, 00185 Roma, Italy; and <sup>6</sup>Laboratoire Matière et Systèmes Complexes, (Centre National de la Recherche Scientifique-Unité Mixte de Recherche 7057), Université Paris VII, 10 rue Alice Domon et Léonie Duquet, 75205 Paris Cedex 13, France

PHYSICAL REVIEW X **6**, 021011 (2016)

## Motility-Driven Glass and Jamming Transitions in Biological Tissues

Dapeng Bi,<sup>1,3</sup> Xingbo Yang,<sup>1,4</sup> M. Cristina Marchetti,<sup>1,2</sup> and M. Lisa Manning<sup>1,2</sup>

Cell motion inside dense tissues governs many biological processes, including embryonic development and cancer metastasis, and recent experiments suggest that these tissues exhibit collective glassy behavior. To make quantitative predictions about glass transitions in tissues, we study a self-propelled Voronoi model that simultaneously captures polarized cell motility and multibody cell-cell interactions in a confluent

Voronoi neighbours

- topological models:  $1^{st}$  or  $2^{nd}$  order flocking transition?
- $2^{nd}$  order arguments  $\left\{ \begin{array}{l} \text{numerics} \rightarrow \text{computationally costly, finite size effects} \\ \text{mean-field hydrodynamics} \rightarrow \text{may be misleading} \end{array} \right.$



# Topological models: a specific transition

- Visual or biological cues  $\rightarrow$  'metric-free' or 'topological' alignment

$k$ -nearest neighbours



## Interaction ruling animal collective behavior depends on topological rather than metric distance: Evidence from a field study

M. Ballerini<sup>1</sup>, N. Cabibbo<sup>2,5</sup>, R. Candelier<sup>1,5</sup>, A. Cavagna<sup>1,6\*</sup>, E. Cisbani<sup>1</sup>, I. Giardina<sup>6</sup>, V. Lecomte<sup>1,7,8,9</sup>, A. Orlandi<sup>1</sup>, G. Parisi<sup>10,11,12\*</sup>, A. Procaccini<sup>1</sup>, and M. Viale<sup>1,5,9</sup>, and V. Zdravkovic<sup>1</sup>

<sup>1</sup>Centre for Statistical Mechanics and Complexity (SMC), Consiglio Nazionale delle Ricerche-Istituto Nazionale per la Fisica della Materia, <sup>2</sup>Dipartimento di Fisica, and <sup>3</sup>Sezione Istituto Nazionale di Fisica Nucleare, Università di Roma "La Sapienza," Piazzale Aldo Moro 2, 00185 Roma, Italy; <sup>4</sup>Istituto Superiore di Sanità, viale Regina Elena 299, 00161 Roma, Italy; <sup>5</sup>Istituto dei Sistemi Complessi (ISC), Consiglio Nazionale delle Ricerche, via dei Taurini 19, 00185 Roma, Italy; and <sup>6</sup>Laboratoire Matière et Systèmes Complexes, (Centre National de la Recherche Scientifique-Unité Mixte de Recherche 7057), Université Paris VII, 10 rue Alice Domon et Léonie Duquet, 75205 Paris Cedex 13, France

PHYSICAL REVIEW X **6**, 021011 (2016)

## Motility-Driven Glass and Jamming Transitions in Biological Tissues

Dapeng Bi,<sup>1,3</sup> Xingbo Yang,<sup>1,4</sup> M. Cristina Marchetti,<sup>1,2</sup> and M. Lisa Manning<sup>1,2</sup>

Cell motion inside dense tissues governs many biological processes, including embryonic development and cancer metastasis, and recent experiments suggest that these tissues exhibit collective glassy behavior. To make quantitative predictions about glass transitions in tissues, we study a self-propelled Voronoi model that simultaneously captures polarized cell motility and multibody cell-cell interactions in a confluent

- topological models:  $1^{st}$  or  $2^{nd}$  order flocking transition?
- $2^{nd}$  order arguments  $\left\{ \begin{array}{l} \text{numerics} \rightarrow \text{computationally costly, finite size effects} \\ \text{mean-field hydrodynamics} \rightarrow \text{may be misleading} \end{array} \right.$
- build a topological field theory  $\rightarrow$  challenging

# The topological Active Ising Model

- Previous full mean-field equation for active Ising

$$\partial_t m = D\nabla^2 m - \nabla(v\rho) + 2\rho\Gamma \sinh\left(\beta\frac{m}{\rho}\right) - 2m\Gamma \cosh\left(\beta\frac{m}{\rho}\right)$$

# The topological Active Ising Model

- Previous full mean-field equation for active Ising

$$\partial_t m = D\nabla^2 m - \nabla(v\rho) + 2\rho\Gamma \sinh\left(\beta\frac{m}{\rho}\right) - 2m\Gamma \cosh\left(\beta\frac{m}{\rho}\right)$$

microscopic flipping  $W_j^\pm = \exp(\pm\beta\frac{m_j}{\rho_j})$

# The topological Active Ising Model

- Previous full mean-field equation for active Ising

$$\partial_t m = D\nabla^2 m - \nabla(v\rho) + 2\rho\Gamma \sinh\left(\beta\frac{m}{\rho}\right) - 2m\Gamma \cosh\left(\beta\frac{m}{\rho}\right)$$

microscopic flipping  $W_j^\pm = \exp(\pm\beta\frac{m_j}{\rho_j})$

local aligning field

# The topological Active Ising Model

- Previous full mean-field equation for active Ising

$$\partial_t m = D\nabla^2 m - \nabla(v\rho) + 2\rho\Gamma \sinh\left(\beta \frac{m}{\rho}\right) - 2m\Gamma \cosh\left(\beta \frac{m}{\rho}\right)$$

microscopic flipping  $W_j^\pm = \exp(\pm\beta \frac{m_j}{\rho_j})$

local aligning field

- Now makes it **topological**  $\rightarrow$  alignment with k-nearest neighbours

# The topological Active Ising Model

- Previous full mean-field equation for active Ising

$$\partial_t m = D\nabla^2 m - \nabla(v\rho) + 2\rho\Gamma \sinh\left(\beta \frac{m}{\rho}\right) - 2m\Gamma \cosh\left(\beta \frac{m}{\rho}\right)$$

microscopic flipping  $W_j^\pm = \exp(\pm\beta \frac{m_j}{\rho_j})$

local aligning field

- Now makes it **topological**  $\rightarrow$  alignment with k-nearest neighbours
- Local interaction range  $y(x)$   $\rightarrow$  adaptation to density fluctuations

# The topological Active Ising Model

- Previous full mean-field equation for active Ising

$$\partial_t m = D\nabla^2 m - \nabla(v\rho) + 2\rho\Gamma \sinh\left(\beta\frac{m}{\rho}\right) - 2m\Gamma \cosh\left(\beta\frac{m}{\rho}\right)$$

microscopic flipping  $W_j^\pm = \exp(\pm\beta\frac{m_j}{\rho_j})$

local aligning field

- Now makes it **topological**  $\rightarrow$  alignment with k-nearest neighbours
- Local interaction range  $y(x)$   $\rightarrow$  adaptation to density fluctuations

- $k = \int_{x-y(x)}^{x+y(x)} \rho(z) dz$

# The topological Active Ising Model

- Previous full mean-field equation for active Ising

$$\partial_t m = D\nabla^2 m - \nabla(v\rho) + 2\rho\Gamma \sinh\left(\beta \frac{m}{\rho}\right) - 2m\Gamma \cosh\left(\beta \frac{m}{\rho}\right)$$

microscopic flipping  $W_j^\pm = \exp(\pm\beta \frac{m_j}{\rho_j})$

local aligning field

- Now makes it **topological**  $\rightarrow$  alignment with  $k$ -nearest neighbours
- Local interaction range  $y(x)$   $\rightarrow$  adaptation to density fluctuations

$$\bullet k = \int_{x-y(x)}^{x+y(x)} \rho(z) dz \quad \bullet \frac{m(x)}{\rho(x)} \rightarrow \tilde{m}(x) = \int_{x-y(x)}^{x+y(x)} \frac{m(z)}{k} dz$$



# The topological Active Ising Model

- Previous full mean-field equation for active Ising

$$\partial_t m = D\nabla^2 m - \nabla(v\rho) + 2\rho\Gamma \sinh\left(\beta \frac{m}{\rho}\right) - 2m\Gamma \cosh\left(\beta \frac{m}{\rho}\right)$$

microscopic flipping  $W_j^\pm = \exp(\pm\beta \frac{m_j}{\rho_j})$

local aligning field

- Now makes it **topological**  $\rightarrow$  alignment with k-nearest neighbours
- Local interaction range  $y(x)$   $\rightarrow$  adaptation to density fluctuations

$$k = \int_{x-y(x)}^{x+y(x)} \rho(z) dz \quad \bullet \quad \frac{m(x)}{\rho(x)} \rightarrow \tilde{m}(x) = \int_{x-y(x)}^{x+y(x)} \frac{m(z)}{k} dz$$



$$\partial_t m = D\nabla^2 m - \nabla(v\rho) + 2\rho\Gamma \sinh(\beta\tilde{m}) - 2m\Gamma \cosh(\beta\tilde{m})$$

# The topological Active Ising Model

- Previous full mean-field equation for active Ising

$$\partial_t m = D\nabla^2 m - \nabla(v\rho) + 2\rho\Gamma \sinh\left(\beta \frac{m}{\rho}\right) - 2m\Gamma \cosh\left(\beta \frac{m}{\rho}\right)$$

microscopic flipping  $W_j^\pm = \exp(\pm\beta \frac{m_j}{\rho_j})$

local aligning field

- Now makes it **topological**  $\rightarrow$  alignment with k-nearest neighbours
- Local interaction range  $y(x)$   $\rightarrow$  adaptation to density fluctuations

$$k = \int_{x-y(x)}^{x+y(x)} \rho(z) dz \quad \bullet \quad \frac{m(x)}{\rho(x)} \rightarrow \tilde{m}(x) = \int_{x-y(x)}^{x+y(x)} \frac{m(z)}{k} dz$$



$$\partial_t m = D\nabla^2 m - \nabla(v\rho) + 2\rho\Gamma \sinh(\beta\tilde{m}) - 2m\Gamma \cosh(\beta\tilde{m})$$

- Linear stability analysis  $\rightarrow$  continuous transition at mean field level

# The topological Active Ising Model

- Previous full mean-field equation for active Ising

$$\partial_t m = D\nabla^2 m - \nabla(v\rho) + 2\rho\Gamma \sinh\left(\beta \frac{m}{\rho}\right) - 2m\Gamma \cosh\left(\beta \frac{m}{\rho}\right)$$

microscopic flipping  $W_j^\pm = \exp(\pm\beta \frac{m_j}{\rho_j})$

local aligning field

- Now makes it **topological**  $\rightarrow$  alignment with k-nearest neighbours
- Local interaction range  $y(x)$   $\rightarrow$  adaptation to density fluctuations

$$k = \int_{x-y(x)}^{x+y(x)} \rho(z) dz \quad \bullet \quad \frac{m(x)}{\rho(x)} \rightarrow \tilde{m}(x) = \int_{x-y(x)}^{x+y(x)} \frac{m(z)}{k} dz$$



$$\partial_t m = D\nabla^2 m - \nabla(v\rho) + 2\rho\Gamma \sinh(\beta\tilde{m}) - 2m\Gamma \cosh(\beta\tilde{m})$$

- Linear stability analysis  $\rightarrow$  continuous transition at mean field level

Protected against fluctuations ?

# The topological Active Ising Model

- Previous full mean-field equation for active Ising

$$\partial_t m = D\nabla^2 m - \nabla(v\rho) + 2\rho\Gamma \underbrace{\sinh\left(\beta \frac{m}{\rho}\right)}_{\text{microscopic flipping}} - 2m\Gamma \underbrace{\cosh\left(\beta \frac{m}{\rho}\right)}_{\text{local aligning field}}$$

$W_j^\pm = \exp(\pm\beta \frac{m_j}{\rho_j})$

- Now makes it **topological**  $\rightarrow$  alignment with k-nearest neighbours
- Local interaction range  $y(x)$   $\rightarrow$  adaptation to density fluctuations

$$\bullet k = \int_{x-y(x)}^{x+y(x)} \rho(z) dz \quad \bullet \frac{m(x)}{\rho(x)} \rightarrow \tilde{m}(x) = \int_{x-y(x)}^{x+y(x)} \frac{m(z)}{k} dz$$



$$\partial_t m = D\nabla^2 m - \nabla(v\rho) + 2\rho\Gamma \sinh(\beta\tilde{m}) - 2m\Gamma \cosh(\beta\tilde{m}) + \sqrt{2\sigma\rho}\eta$$

- Linear stability analysis  $\rightarrow$  continuous transition at mean field level

Protected against **fluctuations** ?

# The fluctuating topological hydrodynamics

- Renormalization of linear Landau term

$$\alpha \longrightarrow \alpha + \frac{\sigma\Gamma}{k} g\left(\beta, \frac{\Gamma k}{v\rho}, \frac{\Gamma D}{v^2}\right) \Rightarrow \alpha \text{ density dependent}$$

# The fluctuating topological hydrodynamics

- Renormalization of linear Landau term

$$\alpha \longrightarrow \alpha + \frac{\sigma\Gamma}{k} g\left(\beta, \frac{\Gamma k}{v\rho}, \frac{\Gamma D}{v^2}\right) \Longrightarrow \alpha \text{ density dependent}$$

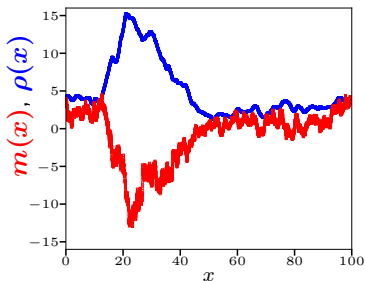
- Linear stability  $\longrightarrow$  discontinuous emergence of flocking

# The fluctuating topological hydrodynamics

- Renormalization of linear Landau term

$$\alpha \longrightarrow \alpha + \frac{\sigma\Gamma}{k} g\left(\beta, \frac{\Gamma k}{v\rho}, \frac{\Gamma D}{v^2}\right) \Longrightarrow \alpha \text{ density dependent}$$

- Linear stability  $\longrightarrow$  discontinuous emergence of flocking
- Simulations of the topological stochastic PDE

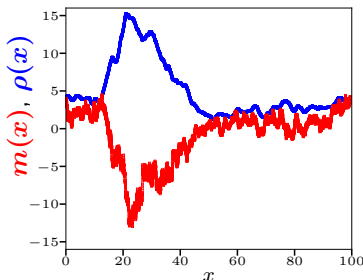


# The fluctuating topological hydrodynamics

- Renormalization of linear Landau term

$$\alpha \longrightarrow \alpha + \frac{\sigma\Gamma}{k} g\left(\beta, \frac{\Gamma k}{v\rho}, \frac{\Gamma D}{v^2}\right) \Longrightarrow \alpha \text{ density dependent}$$

- Linear stability  $\longrightarrow$  discontinuous emergence of flocking
- Simulations of the topological stochastic PDE



- So far  $\longrightarrow$  predictions for field-theoretic models

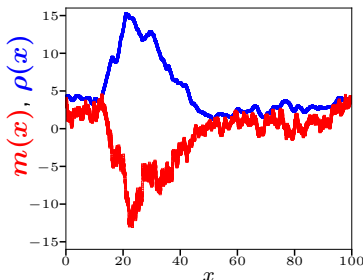


# The fluctuating topological hydrodynamics

- Renormalization of linear Landau term

$$\alpha \longrightarrow \alpha + \frac{\sigma\Gamma}{k} g\left(\beta, \frac{\Gamma k}{v\rho}, \frac{\Gamma D}{v^2}\right) \Longrightarrow \alpha \text{ density dependent}$$

- Linear stability  $\longrightarrow$  discontinuous emergence of flocking
- Simulations of the topological stochastic PDE

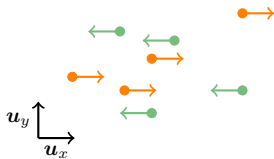


- So far  $\longrightarrow$  predictions for field-theoretic models  
 $\hookrightarrow$  Let's try to see if it is robust for microscopic models !

# The microscopic topological dynamics

- Microscopic dynamics of the **topological** Active Ising Model
  - \* Off-lattice Langevin particles
  - \* Each particles carries a spin

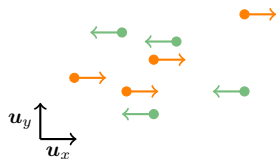
$$\dot{\mathbf{r}}_j = s_j v \mathbf{u}_x + \sqrt{2D} \boldsymbol{\eta}_j$$



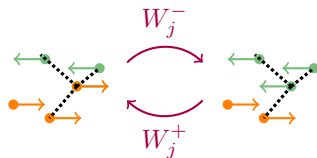
# The microscopic topological dynamics

- Microscopic dynamics of the **topological** Active Ising Model

- \* Off-lattice Langevin particles
- \* Each particles carries a spin



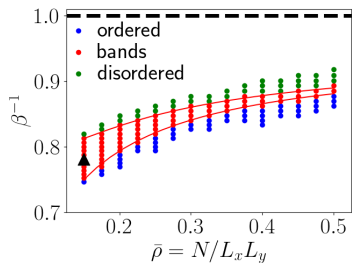
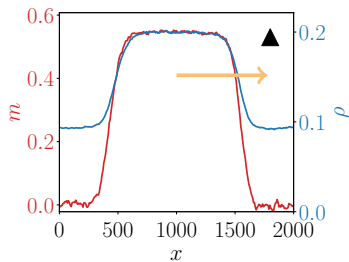
$$\dot{\mathbf{r}}_j = s_j v \mathbf{u}_x + \sqrt{2D} \boldsymbol{\eta}_j$$



- \* Flipping rates  $W_j^\pm = \Gamma \exp(\pm\beta\tilde{m}_j)$  with  $\tilde{m}_j =$  averaged magnetization of  $k$ -nearest neighbours

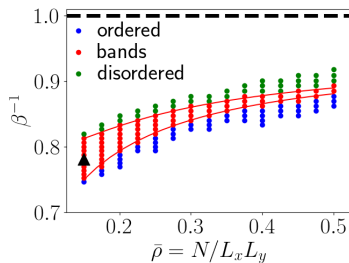
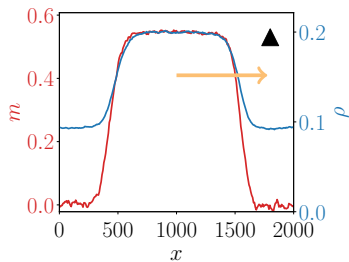
# The microscopic topological dynamics

- Results of the microscopic **topological** Active Ising Model



# The microscopic topological dynamics

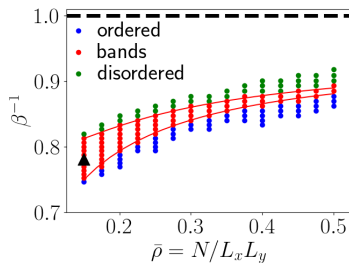
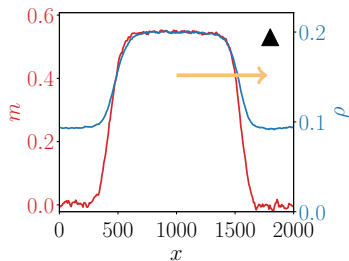
- Results of the microscopic **topological** Active Ising Model



Same fate for microscopic topological models

# The microscopic topological dynamics

- Results of the microscopic **topological** Active Ising Model

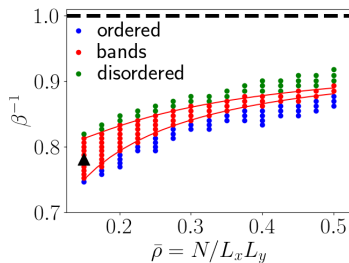
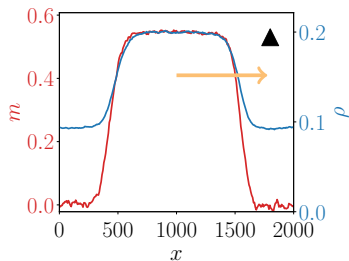


Same fate for microscopic topological models

- Is it **model-dependent** ? Only for active spins ?

# The microscopic topological dynamics

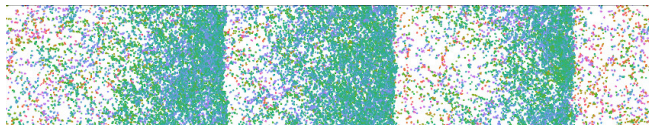
- Results of the microscopic **topological** Active Ising Model



Same fate for microscopic topological models

- Is it **model-dependent** ? Only for active spins ?

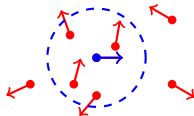
↳ Holds also for **topological Vicsek Model**



# Revisiting the classification

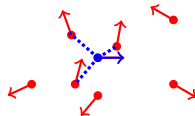
## First order / coexistence

- Vicsek: metrical alignment



## Second order / continuous

- Vicsek: topological alignment



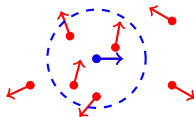
- Active Ising Model: hydrodynamic



# Revisiting the classification

## First order / coexistence

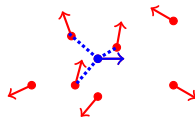
- Vicsek: metrical alignment



- Active Ising Model: numerics

## Second order / continuous

- Vicsek: topological alignment

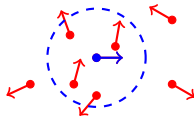


- ~~Active Ising Model: hydrodynamic~~

# Revisiting the classification

## First order / coexistence

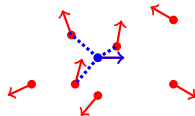
- Vicsek: metrical alignment



- Active Ising Model: numerics

## Second order / continuous

- Vicsek: topological alignment

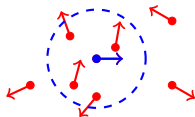


- Active Ising Model: hydrodynamic

# Revisiting the classification

## First order / coexistence

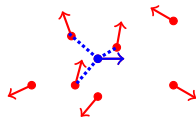
- Vicsek: metrical alignment



- Active Ising Model: numerics

## Second order / continuous

- Vicsek: topological alignment



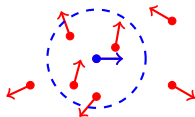
- Active Ising Model: hydrodynamic

- Are all models of collective motion first order ?

# Revisiting the classification

## First order / coexistence

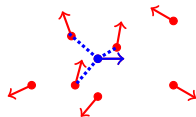
- Vicsek: metrical alignment



- Active Ising Model: numerics

## Second order / continuous

- Vicsek: topological alignment



- Active Ising Model: hydrodynamic

- Fully connected models

- Are all models of collective motion first order ?

↪ No, fully connected alignment  $\Leftrightarrow$  continuous transition

$$\alpha \longrightarrow \alpha + \frac{\sigma\Gamma}{N} g\left(\beta, \frac{\Gamma D}{v^2}\right) \Rightarrow \begin{array}{l} * \text{ no dependence on local density} \\ * \text{ vanishes as the number of particles diverges} \end{array}$$

- Fluctuations renormalize  $T_c$  making it density dependent
- $T_c(\rho)$  turns a **deceptive continuous** transition into a **first order** scenario

- Fluctuations renormalize  $T_c$  making it density dependent
- $T_c(\rho)$  turns a **deceptive continuous** transition into a **first order** scenario
- **Topological** alignment gives no protection  $\rightarrow$  onset of flocking remains **discontinuous**

# Summary and outlook

- Quantification of departure from equilibrium in AOUP
  - ↳ overdamped active particle: go to underdamped scenario
  - ↳ Effect of inertia on nonequilibrium signatures ?

# Summary and outlook

- Quantification of departure from equilibrium in AOUP
  - ↳ overdamped active particle: go to underdamped scenario
  - ↳ Effect of inertia on nonequilibrium signatures ?
- Emergence of MIPS in polar liquid
  - ↳ What about MIPS in flocking bands ? Accessible in experiments ?



# Summary and outlook

- Quantification of departure from equilibrium in AOUP
  - ↳ overdamped active particle: go to underdamped scenario
  - ↳ Effect of inertia on nonequilibrium signatures ?
- Emergence of MIPS in polar liquid
  - ↳ What about MIPS in flocking bands ? Accessible in experiments ?
- $k$ -nearest neighbours alignment discontinuous → generic for other topological rules ?
  - ↳ Necessary and sufficient condition for fluctuation-induced first-order flocking

# Summary and outlook

- Quantification of departure from equilibrium in AOUP
  - ↳ overdamped active particle: go to underdamped scenario
  - ↳ Effect of inertia on nonequilibrium signatures ?
- Emergence of MIPS in polar liquid
  - What about MIPS in flocking bands ? Accessible in experiments ?
- $k$ -nearest neighbours alignment discontinuous → generic for other topological rules ?
  - ↳ Necessary and sufficient condition for fluctuation-induced first-order flocking

- List of publications

[D. Martin](#), J. O'byrne, ME. Cates, É. Fodor, C. Nardini, J. Tailleur, F. Van Wijland, Phys. Rev. E **103**, 032607, (2021)

[D. Martin](#), T. Arnoulx de Pirey, JSTAT Mech. **4**, 043205 (2021)

[D. Geyer](#), [D. Martin](#), J. Tailleur, D. Bartolo, Phys. Rev. X **9**, 031043, (2019)

[D. Martin](#), H. Chaté, C. Nardini, A. Solon, J. Tailleur and F. Van Wijland, Phys. Rev. Lett. **126** 148001 (2021)