

Motility-induced solidification in roller flocks

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Active matter : an out-of-equilibrium field

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- Properties of assemblies of micro-agents dissipating energy in their medium

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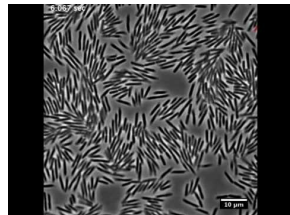
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- Properties of assemblies of micro-agents dissipating energy in their medium
- Canonical examples in nature : bacterial colonies, fish schools, cellular tissues...



A swarm of fish



E. Coli bacteria

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- 2 Focus on Quincke rollers in a racetrack

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- 2 Focus on Quincke rollers in a racetrack
- 3 Motility Induced Phase Separation (MIPS)
- 4 Motility Induced solidification in roller flocks

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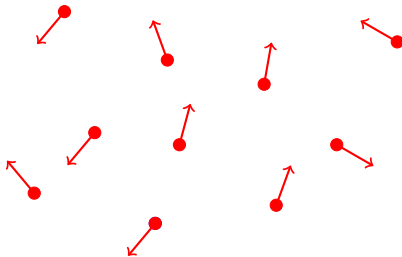
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- The Vicsek model : a canonical model
 - agents moving at constant speed in a given direction
 - noisy aligning interactions between agents



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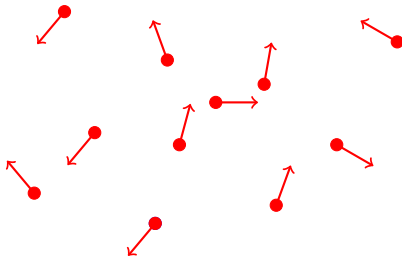
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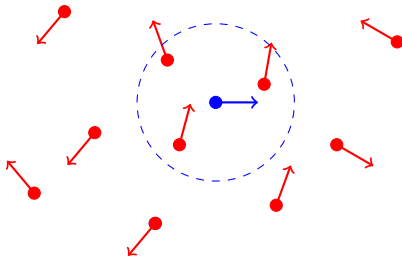
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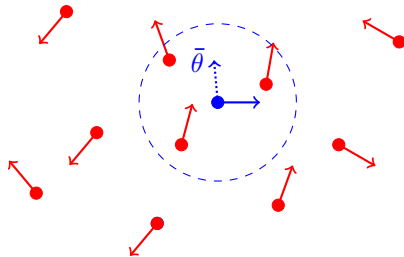
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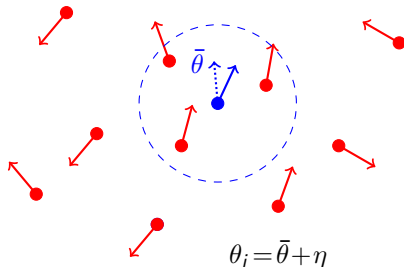
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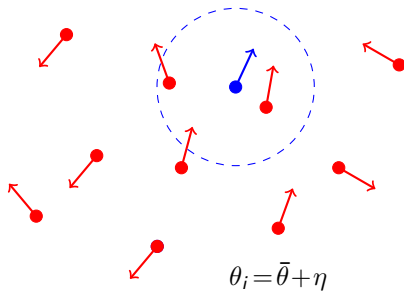
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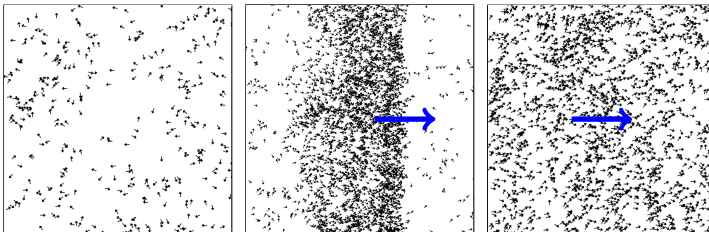
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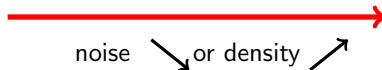
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Disordered

Flocking bands

Polar liquid



- A first order phase transition

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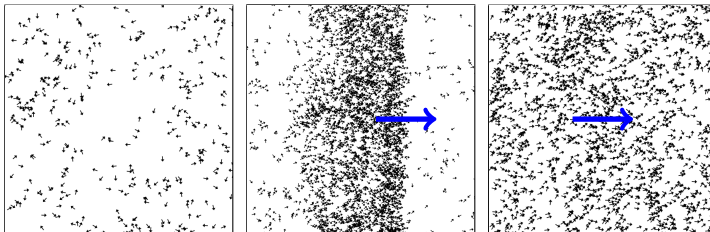
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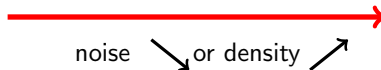
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Disordered

Flocking bands

Polar liquid



- A first order phase transition
- The liquid fraction \nearrow linearly with ρ .

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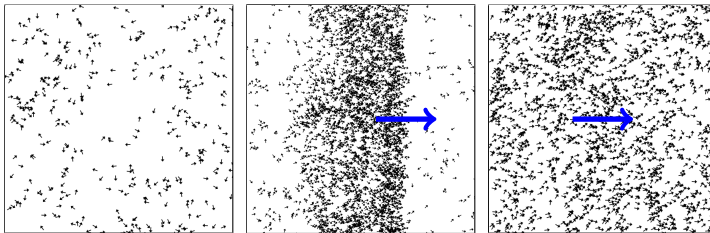
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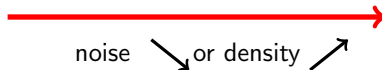
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Disordered

Flocking bands

Polar liquid



- A first order phase transition
- The liquid fraction \nearrow linearly with ρ .
- Hysteresis loop

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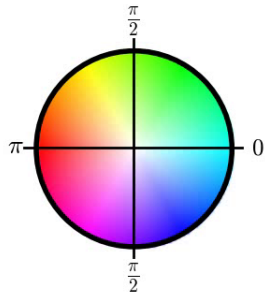
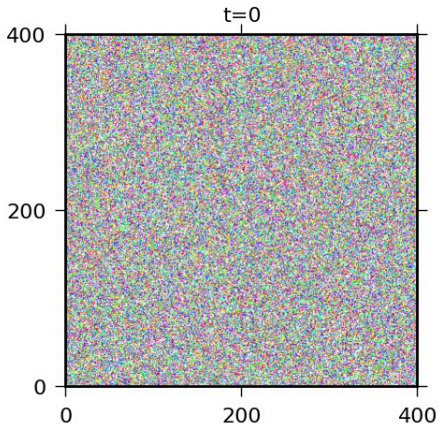
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Toner-Tu theory for Vicsek model

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- state of the art : coarse-grained theories

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- state of the art : coarse-grained theories
- A minimalist 1D Toner-Tu hydrodynamic model for active gas and polar bands coexistence in the literature. [Causin, Solon, . . . , PRL, 2014]

$$\partial_t \rho = D_\rho \partial_{xx} \rho - \partial_x W \quad (1)$$

$$\partial_t W + \lambda W \partial_x W = D_W \partial_{xx} W - \partial_x (v\rho) + a_2 W - a_4 W^3 \quad (2)$$

Toner-Tu theory for Vicsek model

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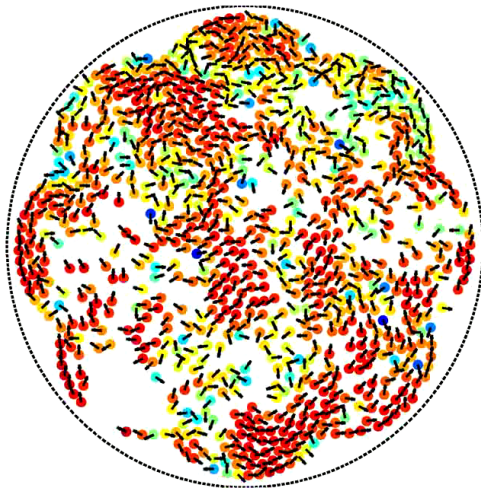
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$$\partial_t W + \lambda W \partial_x W = D_W \partial_{xx} W - \partial_x (v\rho) + a_2 W - a_4 W^3 \quad (2)$$

- Reproduces Vicsek phenomenology
 - Polar bands propagating in a disordered background
 - First order transition : constant binodals, lever rule, hysteresis loop...

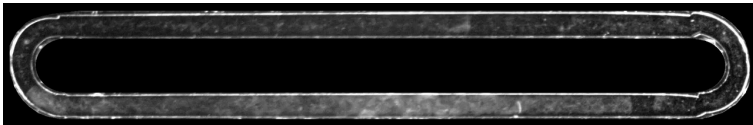
Vicsek phenomenology in experiments

- Shaked grains [Deseigne, Dauchot, Chaté ,PRL, 2010]



Vicsek phenomenology in experiments

- Assemblies of Quincke rollers [Bricard, Caussin, . . . , Nature, 2013]



- We will now focus on this experimental realization of Vicsek model

Physics of Quincke rollers

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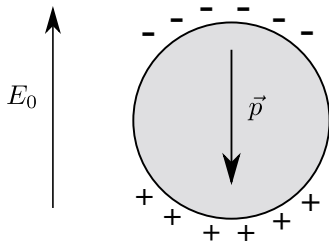
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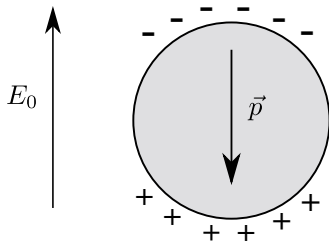
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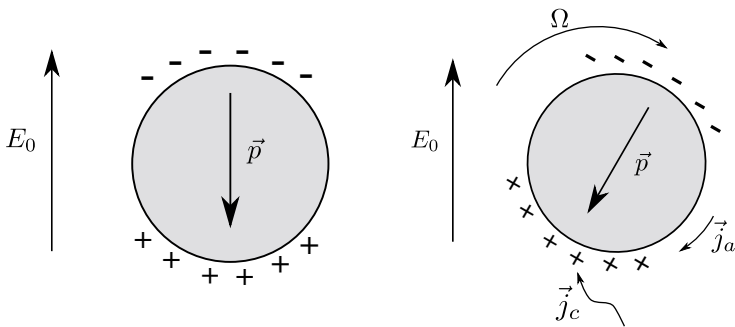
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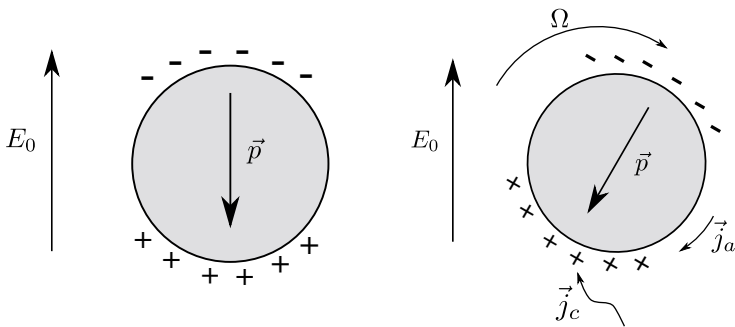
- A roller immersed in a conductive fluid submitted to E_0 develops an electrostatic dipole

Physics of Quincke rollers



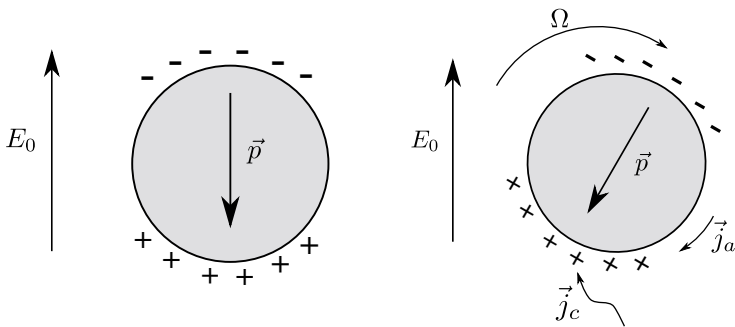
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- Can it sustain steady state rotation ?

Physics of Quincke rollers



- A roller immersed in a conductive fluid submitted to E_0 develops an electrostatic dipole
- Can it sustain steady state rotation ?
- Yes, if on the surface \vec{j}_a compensated by \vec{j}_c
 - \vec{j}_c conductive current generated by \vec{p} and E_0
 - \vec{j}_a advective current due to rotation

Physics of Quincke rollers



- A roller immersed in a conductive fluid submitted to E_0 develops an electrostatic dipole
- Can it sustain steady state rotation ?
- Yes, if on the surface \vec{j}_a compensated by \vec{j}_c
 - \vec{j}_c conductive current generated by \vec{p} and E_0
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- It happens above a threshold field E_Q : $\Omega = \frac{1}{\tau} \sqrt{\left(\frac{E_0}{E_Q}\right)^2 - 1}$

Physics of Quincke rollers

- Rollers experience hydrodynamic and electrostatic interactions

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Physics of Quincke rollers

- Rollers experience hydrodynamic and electrostatic interactions
- Alignment is due to :
 - long range electrostatic dipole-dipole interactions
 - short range hydrodynamic interactions

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- Rollers experience hydrodynamic and electrostatic interactions
- Alignment is due to :
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 - short range hydrodynamic interactions
- Some orders of magnitude :

• radius $a = 5\mu m$	velocity $v = 1mm.s^{-1}$	rotation $\Omega = 1kHz$
• width $l = 2mm$	length $L = 1cm$	height $H = 200\mu m$

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- Dilute coarse-grained evolution quantitatively described by Toner-Tu [Geyer, Morin, Bartolo, Nature Materials, 2018]
- However, missing features at high density

Beyond Vicsek phenomenology at high density

- Unveiling new phase transition at high ρ for assemblies of Quincke rollers.

Beyond Vicsek phenomenology at high density

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- Coexistence of polar liquid and counter-propagating traffic jams



Beyond Vicsek phenomenology at high density

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- We call it active solidification

Beyond Vicsek phenomenology at high density

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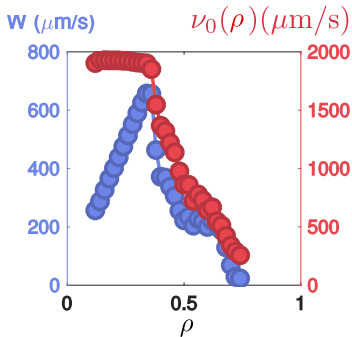
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- What is happening ?

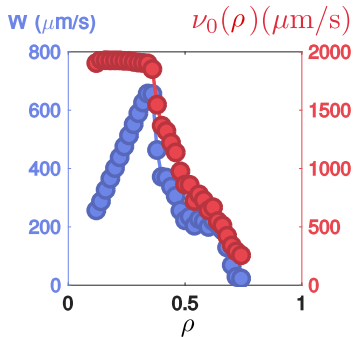
Beyond Vicsek phenomenology at high density

- What is happening ?



Beyond Vicsek phenomenology at high density

- What is happening ?



- Drop of velocity and polar order as $\rho \nearrow$

Beyond Vicsek phenomenology : MIPS ?

- $v(\rho) \searrow$ when $\rho \nearrow$: MIPS ingredient

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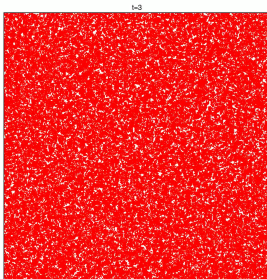
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Beyond Vicsek phenomenology : MIPS ?

- $v(\rho) \searrow$ when $\rho \nearrow$: MIPS ingredient
- Self-propelled particles with pairwise forces (PFAPs)
[Fily & Marchetti PRL 2012, Redner et al. PRL 2013, Stenhammar et al. PRL 2013, Bialké et al. PRL 2013, ...]

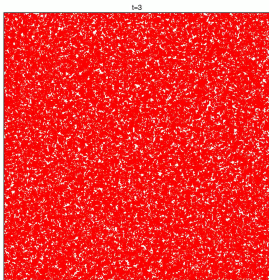
$$\dot{\mathbf{r}}_i = v\mathbf{u}(\theta_i) + \mu \sum_j F_{ij}(\mathbf{r}_i - \mathbf{r}_j) + \sqrt{2D_t}\eta_i; \quad \dot{\theta}_i = \sqrt{2D_r}\xi_i$$



Beyond Vicsek phenomenology : MIPS ?

- $v(\rho) \searrow$ when $\rho \nearrow$: MIPS ingredient
- **Self-propelled particles with pairwise forces (PFAPs)**
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$$\dot{\mathbf{r}}_i = v\mathbf{u}(\theta_i) + \mu \sum_j F_{ij}(\mathbf{r}_i - \mathbf{r}_j) + \sqrt{2D_t}\eta_{it}; \quad \dot{\theta}_i = \sqrt{2D_r}\xi_i$$



- Interactions yields **decreasing** $v(\rho) \equiv \sum_i \dot{\mathbf{r}}_i \cdot \vec{u}(\theta_i)$ [Fily et al PRL (2012)]

MIPS scenario

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- Non-uniform speed $\rightarrow \partial_t P = -\nabla(v(\mathbf{r})\mathbf{u}(\theta)P) + \Theta P$

Accumulation in slow regions $\rho \sim \frac{1}{v(\mathbf{r})}$

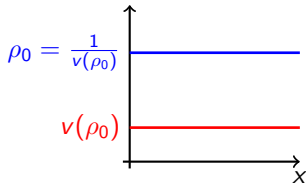
- $v'(\rho) < 0 \rightarrow$ Slow down in dense regions

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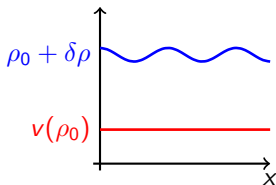


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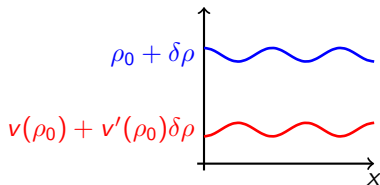


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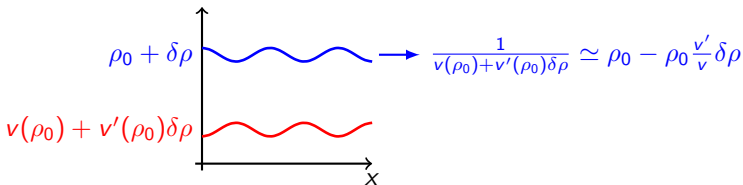


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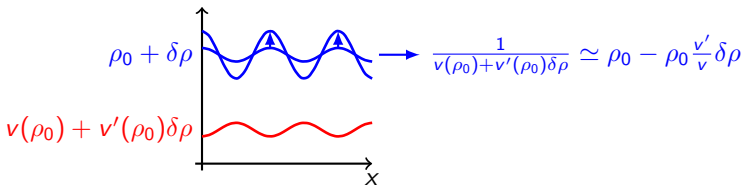


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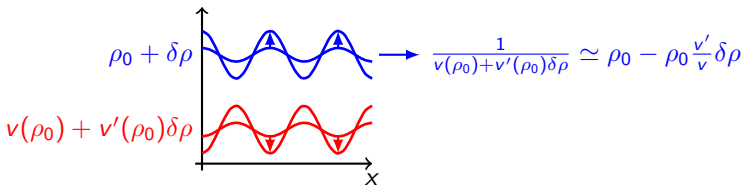
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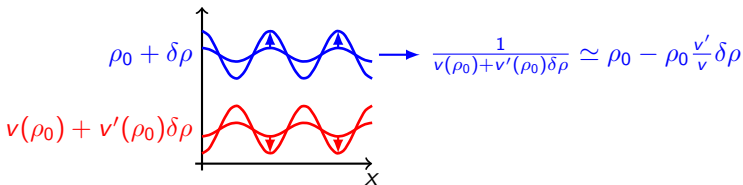
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- $v'(\rho) < 0 \rightarrow$ Slow down in dense regions



- Linear instab. if $\rho_0 v' + v \leq 0$
- What does it entail ? \rightarrow Motility Induced Phase Separation

Including MIPS in Toner-Tu hydrodynamics

- Take again the Toner-Tu hydrodynamic for Vicsek

$$\partial_t \rho = D_\rho \partial_{xx} \rho - \partial_x W \quad (3)$$

$$\partial_t W + \lambda W \partial_x W = D_W \partial_{xx} W - \partial_x (v\rho) + a_2 W - a_4 W^3 \quad (4)$$

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- Rollers' velocity drop at high density : $v \rightarrow v(\rho)$
- Rollers lose orientational order at high density : $a_2 \rightarrow a_2(\rho)$

Including MIPS in Toner-Tu hydrodynamics

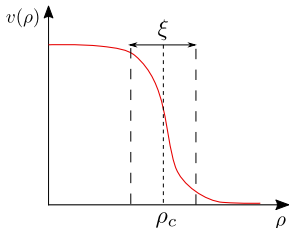
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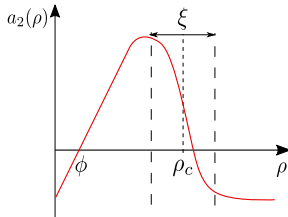
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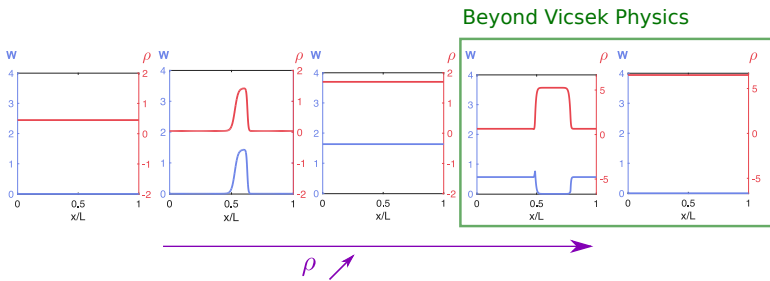
Velocity drop



Ordering drop

Modified Toner-Tu hydrodynamics

- **New phase transition**



Phase diagram of modified Toner-Tu

- Linear stability exhibits two MIPS-like criteria
 - **solid melting** : $v(\rho) + \rho v'(\rho) < -K_1 D_\rho$
 - **solid nucleation** : $v(\rho) + \rho v'(\rho) < -K_2 W_0^2$

Modified Toner-Tu transition is first order

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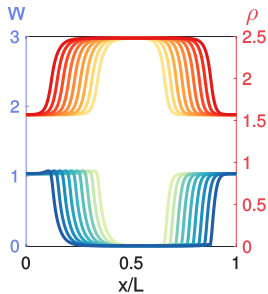
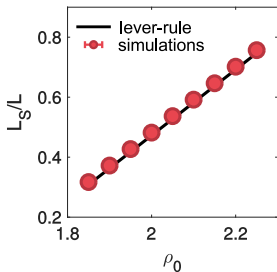
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- Finite lower bound for traffic jam extent and lever rule

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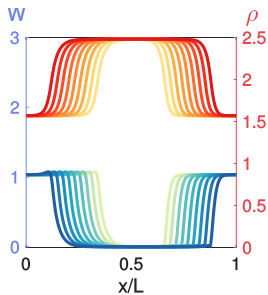
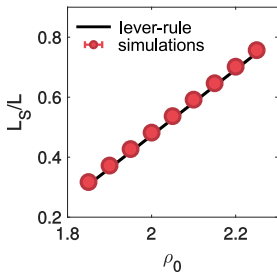
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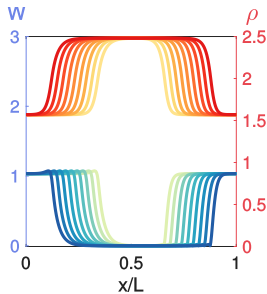
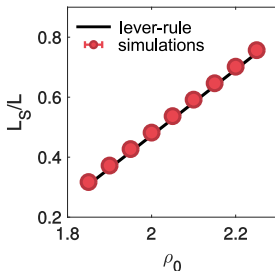
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- Finite lower bound for traffic jam extent and lever rule
- Constants binodals at coexistence

Modified Toner-Tu transition is first order



- Finite lower bound for traffic jam extent and lever rule
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Modified Toner-Tu transition is first order

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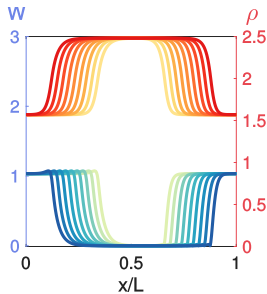
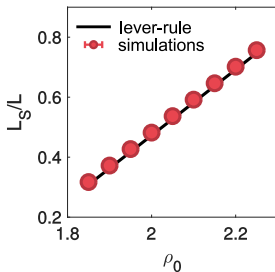
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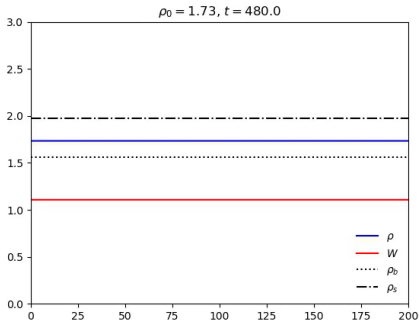
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- Back to experiments \rightarrow same phenomenology ?

Phase diagram of Quincke rollers

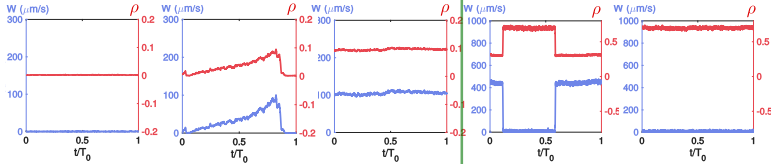
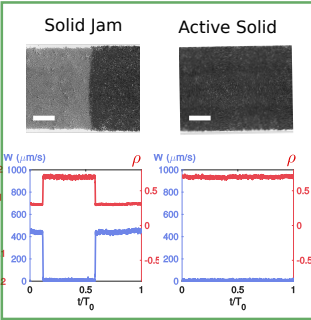
Motility-induced solidification in roller flocks

D. Martin

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Beyond Vicsek Physics



Phase diagram of Quincke rollers

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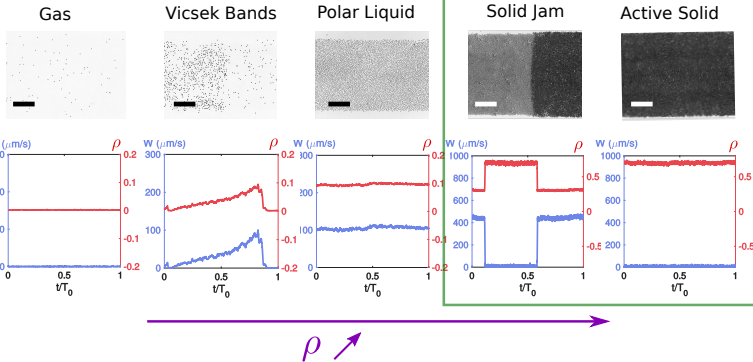
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- Same phase diagram

Active solidification is a first order transition

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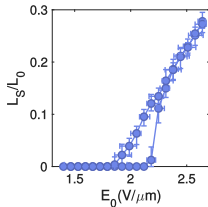
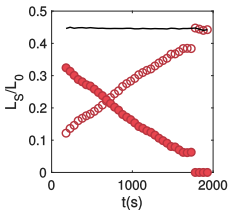
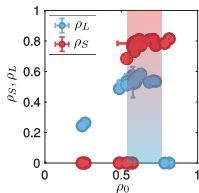
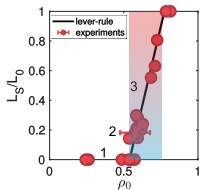
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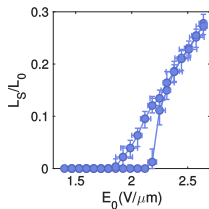
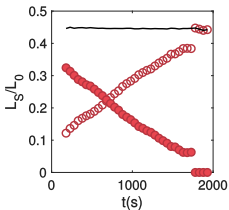
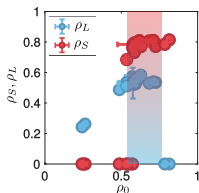
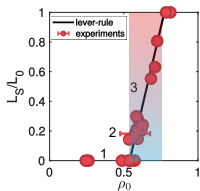
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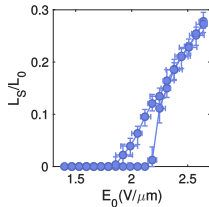
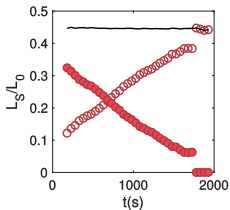
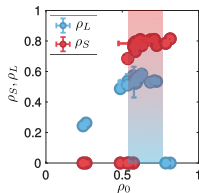
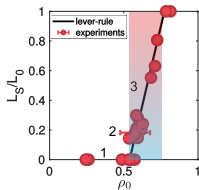


Active solidification is a first order transition



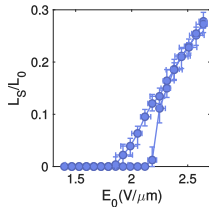
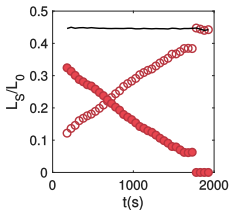
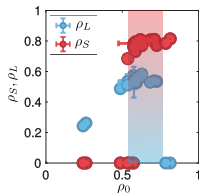
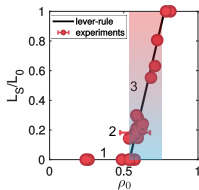
- Finite lower bound for traffic jam extent and lever rule

Active solidification is a first order transition



- Finite lower bound for traffic jam extent and lever rule
- Constant binodals at coexistence

Active solidification is a first order transition



- Finite lower bound for traffic jam extent and lever rule
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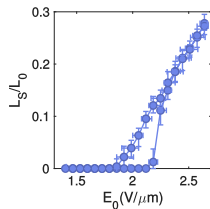
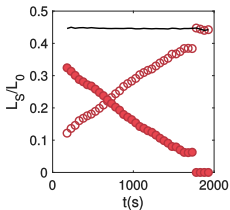
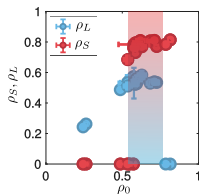
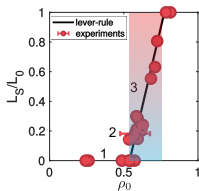
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Conclusion and outlook

- Unveiling of a new phase transition at high ρ in roller flock : active solidification

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Conclusion and outlook

- Unveiling of a new phase transition at high ρ in roller flock : **active solidification**
- It can be described by **MIPS** occurring in a polar flock
- It is a generic feature of aligning motile polar units and speed reduction at high density
- A lot of open questions remaining :
 - Is there other phase transition to discover at high density ?
 - What is the dynamic and the structure of the jammed phase ?

